In search of topological states with half quantum vortices.

Eun-Ah Kim Cornell University

Suk-Bum Chung (Stanford) Hendrik Bluhm (Harvard) Subroto Mukerjee (Berkeley) Daniel Agterberg (UWM)









YBCO single crystal @ 22.3K. Scan height 420nm



In search of topological states with half quantum vortices

- Topological order and fractionalization
- 1/2 QV's
- Stability of 1/2-QV's in SrRuO
- 1/2 QV lattices



Conventional order

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 Symmetry of the underlying Hamiltonian.

Conventional order



Symmetry of the underlying Hamiltonian.
 reduced symmetry

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Local measurements
 order parameter

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Gapped spectrum
 Topological invariance
 mergent symmetry

No local order parameter.

Topological degeneracy N_g.

Sweet Topology



Sweet Topology



Sweet Topology



Fractional charge e*=e/q Ng=q^g e.g., N1=3



Wen and Niu, PRB, 1990 Stone and Chung, PRB, 2006

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2n Non-abelian vortices

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2n Non-abelian vortices



 $N_{2n}=2^{n-1}$ for MR state or p+ip SF

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©n- <u>abelian</u> vortex states

In <u>abelian</u> vortex states

 $\Psi(x_1, \cdots, x_n) = c$ -number

On- <u>abelian</u> vortex states

 $\Psi(x_1, \cdots, x_n) = c$ -number exchange of qp's: phase multiplication to a complex number

 $\Psi(x_1 \leftrightarrow x_3) = e^{i\theta}\Psi$ $\Psi(x_1 \leftrightarrow x_2) = e^{i\theta}\Psi$

In a nonabelian vortex states ⇒set of Qubits

In nonabelian vortex states set of Qubits



In nonabelian vortex states set of Qubits


Nonabelian statistics

In nonabelian vortex states set of Qubits



 $\Psi(x_1, \cdots, x_n) = \begin{pmatrix} \psi_1 \\ \vdots \\ \psi_{d(n)} \end{pmatrix} \longrightarrow \begin{array}{l} \text{exchange of qp's:} \\ \text{rotation in } d(n) \text{ dim} \\ \text{Hilbert space} \end{array}$

Nonabelian statistics

In nonabelian vortex states set of Qubits



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 $\Psi(x_1 \leftrightarrow x_3) = M\Psi(x_1, \cdots, x_n)$ $\Psi(x_1 \leftrightarrow x_2) = N\Psi(x_1, \cdots, x_n)$

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A. Leggett, RMP 1975 Sigrist & Ueda RMP 1991

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$$\oslash$$
 Gap function $\Delta_{ss'}(\mathbf{k}) = -\sum_{\mathbf{k}', s_3, s_4} V_{s'ss_3s_4}(\mathbf{k}, \mathbf{k}') \langle a_{\mathbf{k}'s_3} a_{-\mathbf{k}'s_4} \rangle$

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$$\widehat{\Delta}(\mathbf{k}) = -\widehat{\Delta}^{T}(-\mathbf{k})$$

Singlet gap function

a

$$\widehat{\Delta}(\mathbf{k}) = i \widehat{\sigma}_{y} \psi(\mathbf{k}) = \begin{bmatrix} 0 & \psi(\mathbf{k}) \\ -\psi(\mathbf{k}) & 0 \end{bmatrix}$$

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Triplet gap matrix

$$\widehat{\Delta}(\mathbf{k}) = i(\mathbf{d}(\mathbf{k}) \cdot \widehat{\boldsymbol{\sigma}}) \widehat{\boldsymbol{\sigma}}_{y}$$

$$= \begin{bmatrix} -d_{x}(\mathbf{k}) + id_{y}(\mathbf{k}) & d_{z}(\mathbf{k}) \\ d_{z}(\mathbf{k}) & d_{x}(\mathbf{k}) + id_{y}(\mathbf{k}) \end{bmatrix}$$

p+ip SC

T-breaking (ABM) $\hat{\Delta}(\mathbf{k}) = \Delta_0(T)(p_x \pm ip_y) \begin{pmatrix} -d_x + id_y & d_z \\ d_z & d_x + id_y \end{pmatrix}$ where $\hat{\mathbf{d}}$ is a real unit vector 🔊 In plane d $d_z=0$ i.e., $\mathbf{d} = (\cos\alpha, \sin\alpha, 0)$ $\Delta(\mathbf{\hat{k}}) = \Delta_0(T)(p_x \pm ip_y) \begin{pmatrix} -e^{-i\alpha} & 0\\ 0 & e^{i\alpha} \end{pmatrix}$

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I/2 QV when $\mathbf{d} = (\cos \alpha, \sin \alpha, \theta)$

The gap matrix $\hat{\Delta}(\mathbf{k}) = \Delta_0(T)(p_x \pm ip_y) \begin{pmatrix} -e^{-i\alpha} & 0\\ 0 & e^{i\alpha} \end{pmatrix}$

I/2 QV when d = (cos α , sin α , 0) 2π winding for only one spin component

The gap matrix

Δ(̂k) = Δ₀(T)(p_x ± ip_y) (-e^{-iα} 0
0 e^{iα})

1/2 QV when d = (cosα, sinα, 0)
2π winding for only one spin component
π winding of order parameter phase φ

+ π rotation of **d** vector



Overtices of p+ip SF zero modes at the core

Vortices of p+ip SF > zero modes at the core Kopnin and Salomaa PRB (1991)

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Zero modes are Majorana

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Zero modes are Majorana

BdG qp's \(\gamma_i^{\dagger} = u\psi_i^{\dagger} + v\psi_i \) \(\gamma_i^{\dagger}(E_n) = \(\gamma_i(-E_n))\) \(\zeta_i) = \(\gamma_i(0)) = \(\gamma_i(0))\) \(\zeta_i) = \(\zeta_i) = \(\zeta_i)\) \(\zeta_i) = \(\zeta_i) = \(\zeta_i)\) \(\zeta_i) = \(\zeta_i) = \(\zeta_i)\) \(\zeta_i) = \(\zeta_i) = \(\zeta_i) = \(\zeta_i)\) \(\zeta_i) = \(\zeta_i) = \(\zeta_i) = \(\zeta_i)\) \(\zeta_i) = \(\zeta_i) = \(\zeta_i) = \(\zeta_i) = \(\zeta_i)\) \(\zeta_i) = \(\zeta_i) = \(\zeta_i) = \(\zeta_i)\) \(\zeta_i)\) \(\zeta_i) = \(\zeta_i)\) \(\zeta_i)\) \(\zeta_i)\) \(\zeta_i) = \(\zeta_i)\) \(\zeta_i)\)

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zero mode: \(\gamma_i^{\dagger}(0) = \(\gamma_i(0)\) \)

Das Sarma, Tewari, Nayak (06) Stone & Chung(06) Ivanov(01)

 $\gamma \gamma \gamma$ $\gamma \gamma$

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Majorana + vortex composite

non-Abelian statistics



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1/2 QV's: single Majorana zero mode

5/2 state described as p+ip paired stated of composite fermion

Pfaffian is real space many body BCS wave function of p+ip SF

HQV is equivalent to 1/4 qp

Moore & Read (91) Read & Green (00) Schriffer, p 48

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K. Ishida et al, Nature (1998)

Spin-triplet superconductivity in Sr₂RuO₄ identified by ¹⁷O Knight shift





Experiments?

NMR on ³He-A thin films: X Hakonen et al. Physica (89)

Small angle neutron scattering: X Riseman et al. Nature (98)

Scanning SQUID imaging: X Dolocan et al, PRL (05), Bjorsson et al, PRB (05)

So NMR in the presence of H ⊥ ab
I d // ab: for H_⊥ ≈ 200 G , Murakawa et al, PRL (04)



\bigcirc Energy competition between full-QV and 1/2-QV

The energy competition between full-QV and 1/2-QV

Reducing vorticity saves magnetic energy

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 d-vector bending costs energy

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 \bigcirc Gradient free energy when $d \perp L$ (London limit)

$$f_{\rm grad}^{\rm 2D} = \frac{1}{2} \left(\frac{\hbar}{2m} \right)^2 \left[\rho_{\rm s} \left(\nabla_{\!\!\perp} \phi - \frac{2e}{\hbar c} \mathbf{A} \right)^2 \! + \rho_{\rm sp} \left(\nabla_{\!\!\perp} \alpha \right)^2 \right] \! + \! \frac{1}{8\pi} (\nabla \times \mathbf{A})^2$$

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Spin current energy diverges logarithmically!

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Spin current energy diverges logarithmically!

$$\epsilon_{\rm sp} = \frac{\pi}{4} \left(\frac{\hbar}{2m}\right)^2 \rho_{\rm sp} \ln\left(\frac{R}{\xi}\right)$$

stability of 1/2 QV



$$\begin{split} E_{\text{pair}}^{\text{half}}(r_{12}) &= \frac{1}{2} \frac{\Phi_0^2}{16\pi^2 \lambda^2} \bigg[\ln \bigg(\frac{\lambda}{\xi} \bigg) + K_0 \bigg(\frac{r_{12}}{\lambda} \bigg) + \frac{\rho_{\text{sp}}}{\rho_{\text{s}}} \ln \bigg(\frac{r_{12}}{\xi} \bigg) \bigg] \\ E^{\text{full}} &= \pi \bigg(\frac{\hbar}{2m} \bigg)^2 \rho_{\text{s}} \ln \bigg(\frac{\lambda}{\xi} \bigg) = \frac{\Phi_0^2}{16\pi^2 \lambda^2} \ln \bigg(\frac{\lambda}{\xi} \bigg) \end{split}$$

Competition between screened magnetic repulsion and unscreened spin attraction
E^{1/2qy-pair}(r)-E^{full-qy}

Finite equilibrium size
for small ρ_{sp}/ρ_s

Leggett RMP 75



Mesoscopic sample

Sample of size ~ λ a few micron



 $L = 2\lambda$





Underway in Budakian lab

Sample of size ~ λ a few micron
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Natural way to stabilize 1/2 QV

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- O Potential of tuning ρ_{sp}/ρ_s

-Knowledge exist for ρ_{sp}/ρ_s as a function of Fermi liquid parameters -p-wave Feshbach resonance

SC (SF) with additional U(1) symmetry due to d rotation

Interlacing lattices of two types of vortices





-Different geometry depending on density and LL mixing

Similar case arise in spinor condensate



Muller & Ho(02) Barnett, Mukergee & Moore(08)

Prediction

Minimze GL free energy to determine the VL structure

Quartic terms in the free energy determine the structure



Field distribution as can be measured by neutron

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Field distribution as can be measured by neutron

Stiffness engineering?

p-wave Feshbach resonance can allow for tuning for Fermi liquid parameters

$$H = \sum_{\mathbf{p}} \epsilon(\mathbf{p}) a_{\mathbf{p}}^{\dagger} a_{\mathbf{p}} + \sum_{\mathbf{p},\alpha} \left[\epsilon_{\alpha} + \frac{\epsilon(\mathbf{p})}{2} \right] b_{\mathbf{p}\alpha}^{\dagger} b_{\mathbf{p}\alpha} + \frac{1}{\sqrt{V}} \sum_{\mathbf{p},\mathbf{q},\alpha} g_{\mathbf{p}} p_{\alpha} \left(b_{\mathbf{q}\alpha} a_{\mathbf{p}+\frac{\mathbf{q}}{2}}^{\dagger} a_{-\mathbf{p}+\frac{\mathbf{q}}{2}}^{\dagger} + \text{h.c.} \right)$$

Gurarie, L. Radzihovsky, & A. V. Andreev (05)

Observe to arrive at a PD where ρ_{sp}/ρ_s can be tuned as a function of microscopic parameter





@1/2 QV's are not stable in bulk systems



1/2 QV's are not stable in bulk systems
Mesoscopic samples could favor 1/2 QV's



1/2 QV's are not stable in bulk systems
Mesoscopic samples could favor 1/2 QV's
1/2 QV Vortex Lattice can be pursued and detected