

Investigating Magnetic Order in Mesoscopic Superconductors Using Cantilever Torque Magnetometry



Raffi Budakian
Joonho Jang



1 μm

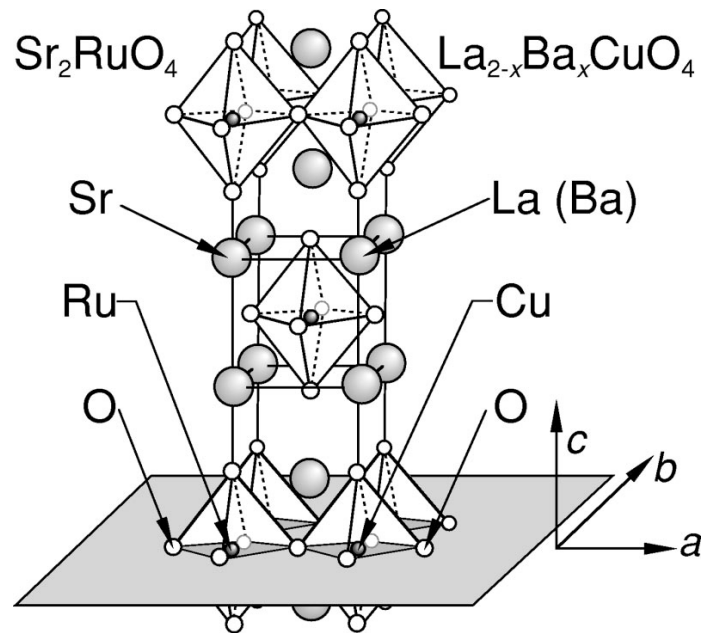
Outline

- Introduction to superconductivity in Sr_2RuO_4
- Describe torque magnetometry measurements to detect moments from edge currents in mesoscopic SRO samples
- Measurements in NbSe_2 (Model system that is layered and s-wave)

SRO Measurements

- Preliminary evidence of moment due to edge currents
- Nonlinear diamagnetic susceptibility
- Concluding Remarks

Superconductivity in Sr_2RuO_4



- Layered perovskite structure similar to high T_c cupra
- Normal state is metallic
- $T_c = 1.5 \text{ K}$

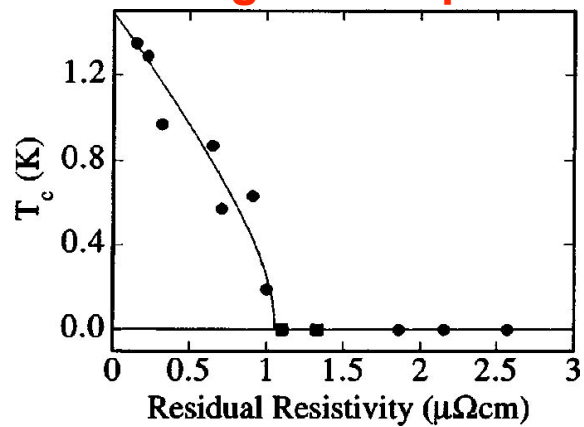
$$a, b = 0.38 \text{ nm}$$

$$c = 1.27 \text{ nm}$$



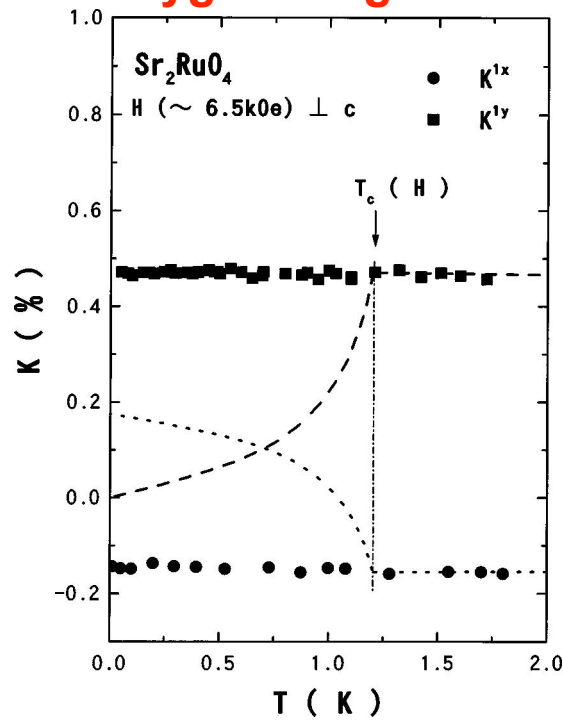
Evidence for Unconventional Superconductivity

Suppression of T_c from non-magnetic impurities



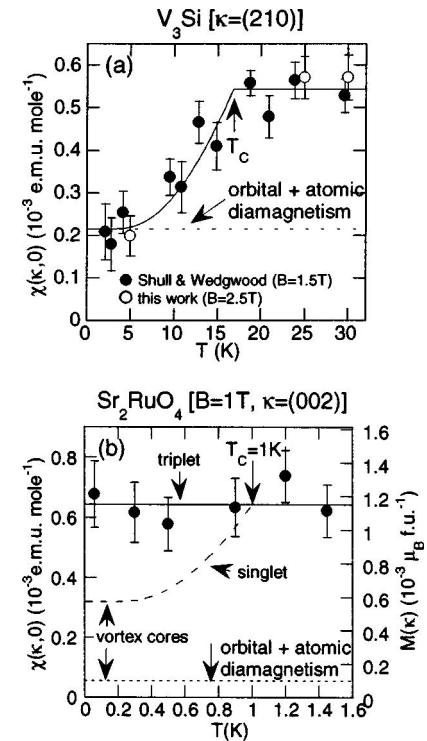
Mackenzie *et al.* (1998)

(NMR) Oxygen Knight shift



Ishida *et al.* (1998)

Spin-polarized neutron scattering



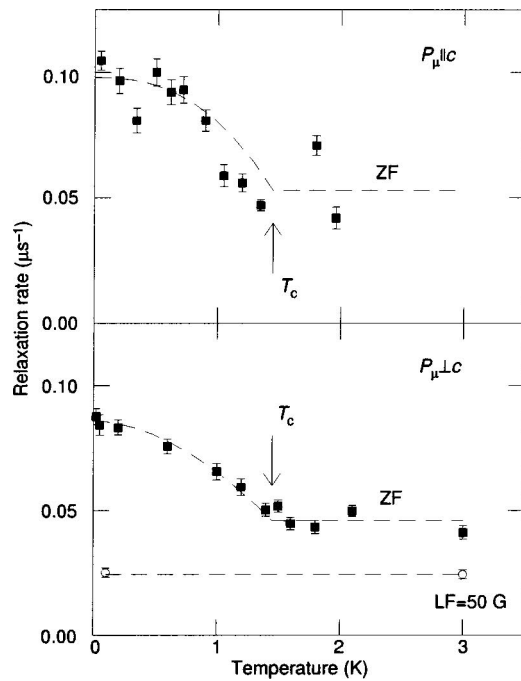
Duffy *et al.* (2000)

$$\vec{S} = 1$$

Spin component of wavefunction is even parity

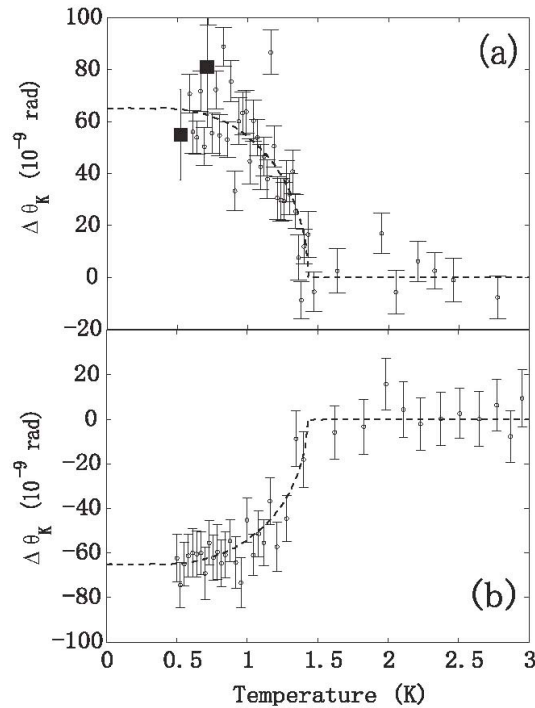
Evidence for Time Reversal Symmetry Breaking

μ SR



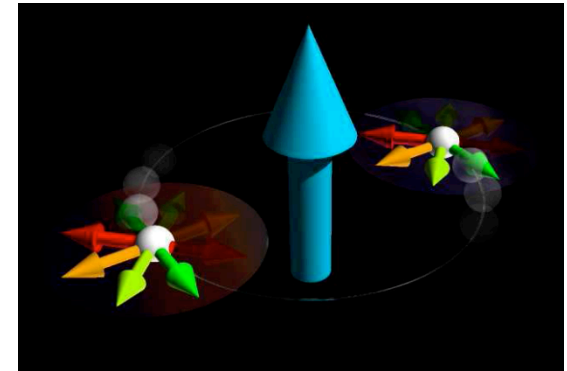
Luke *et al.* (1998)

Kerr Rotation



Xia *et al.* (2006)

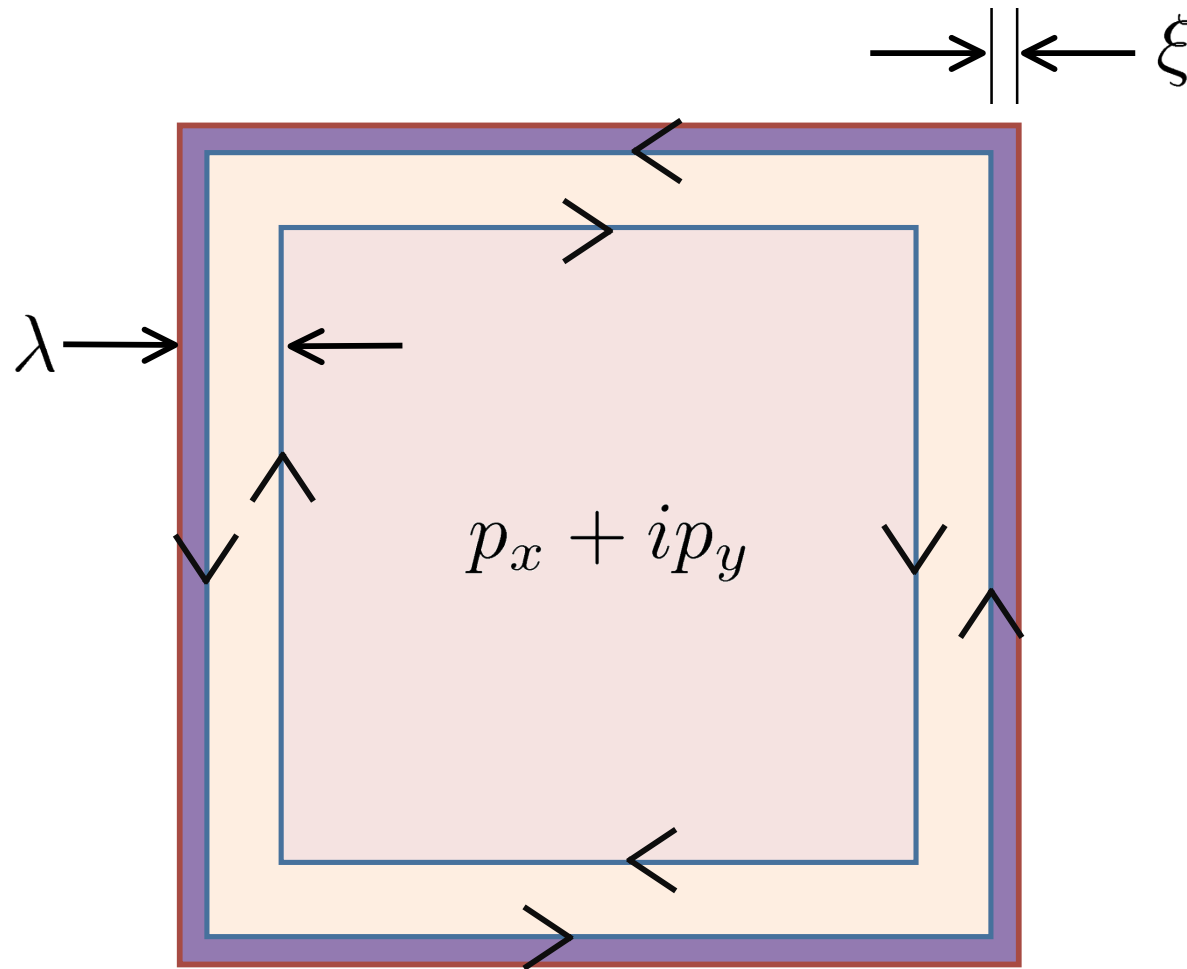
$p_x \pm ip_y$
Complex order parameter








$$\vec{L} = 1, \vec{S} = 1$$

- Domains with orbital order have a net magnetic moment.
- Magnetic fields from domains are screened by the collective motion of CP.

Screening Currents Around Chiral Domains

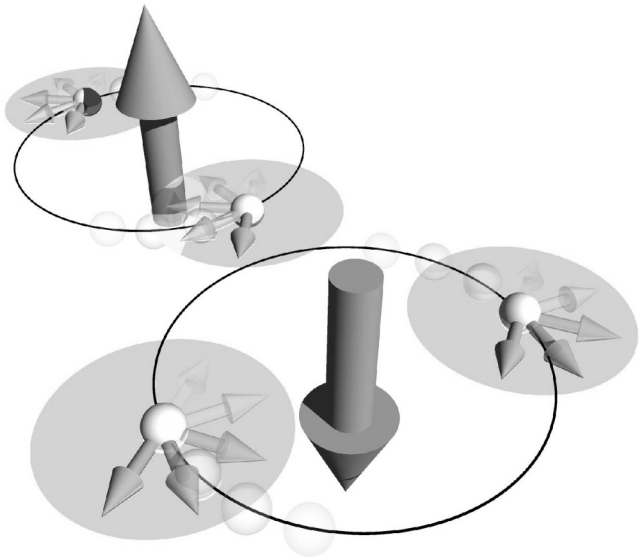


Experimental Evidence for TRS Breaking

Experiment	Status	Domain Size
μ SR		-----
Josephson tunneling		$< 1 \mu\text{m}$
Kerr Rotation		$\sim 50 - 100 \mu\text{m}$
SQUID Imaging		$< 2 \mu\text{m}$
Hall probe Imaging		$< 1 \mu\text{m}$

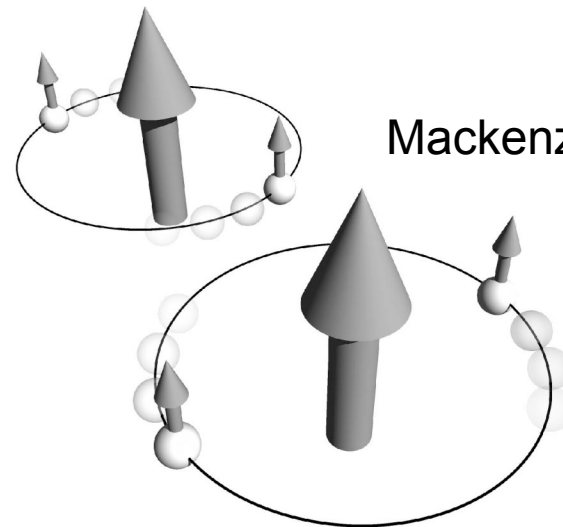
Vortex Matter in Chiral Superconductors

- 2 Vectors are needed to describe the SC order
 - (1) **d**-vector - direction normal to the spin polarization
 - (2) **L**-vector - the angular momentum



$$\vec{d} \parallel \vec{L}$$

Integer quantum vortex:
Orbital phase winds by 2π
d-vector is stationary



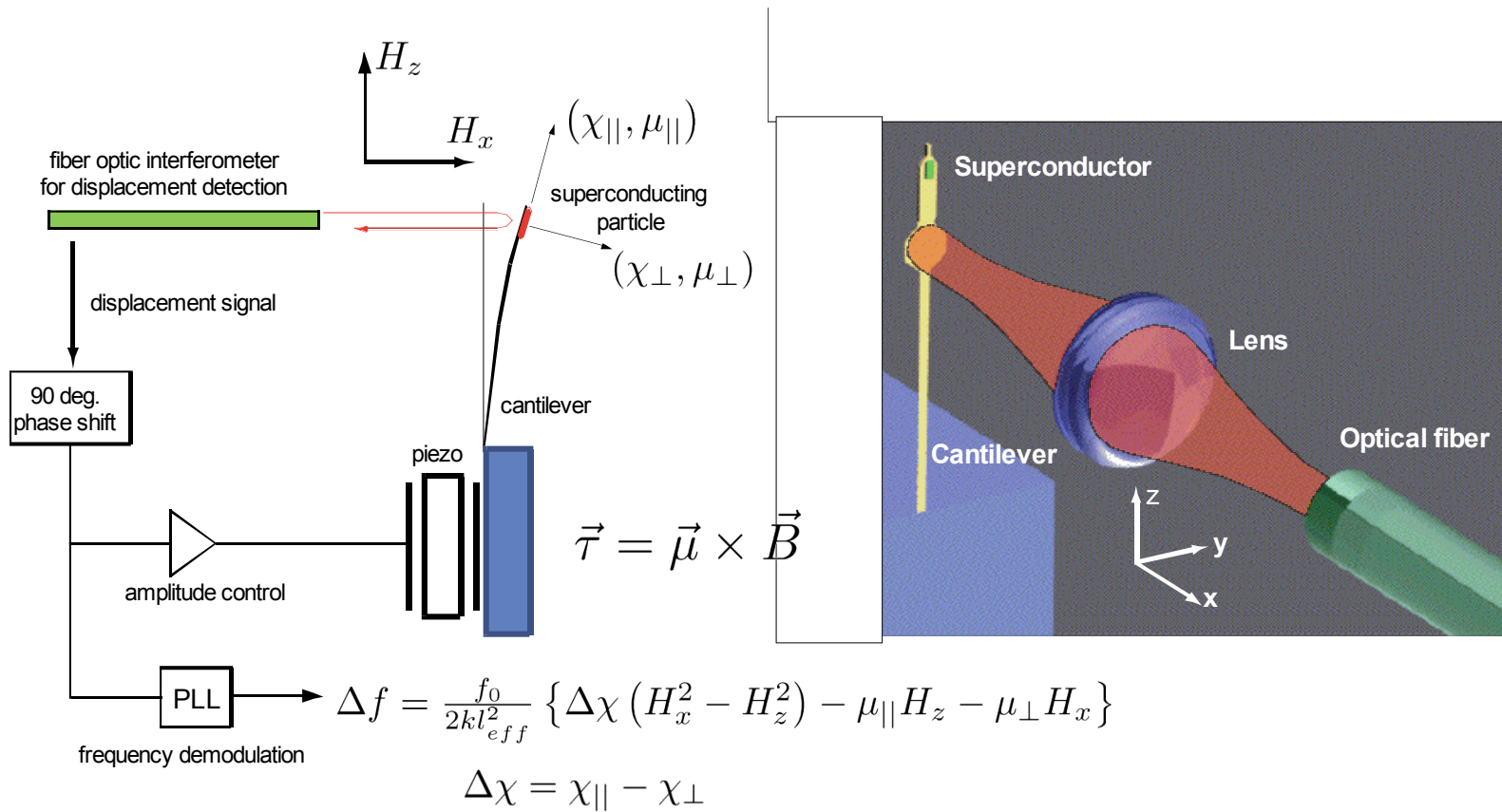
Mackenzie (2003)

$$\vec{d} \perp \vec{L}$$

Half-integer quantum vortex:
Orbital phase winds by π
d-vector winds by $\pm \pi$



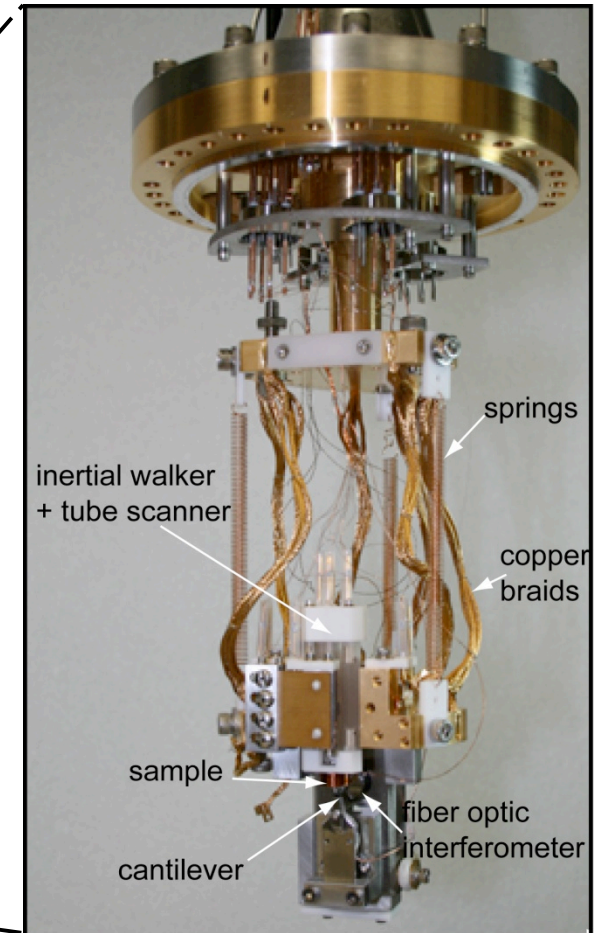
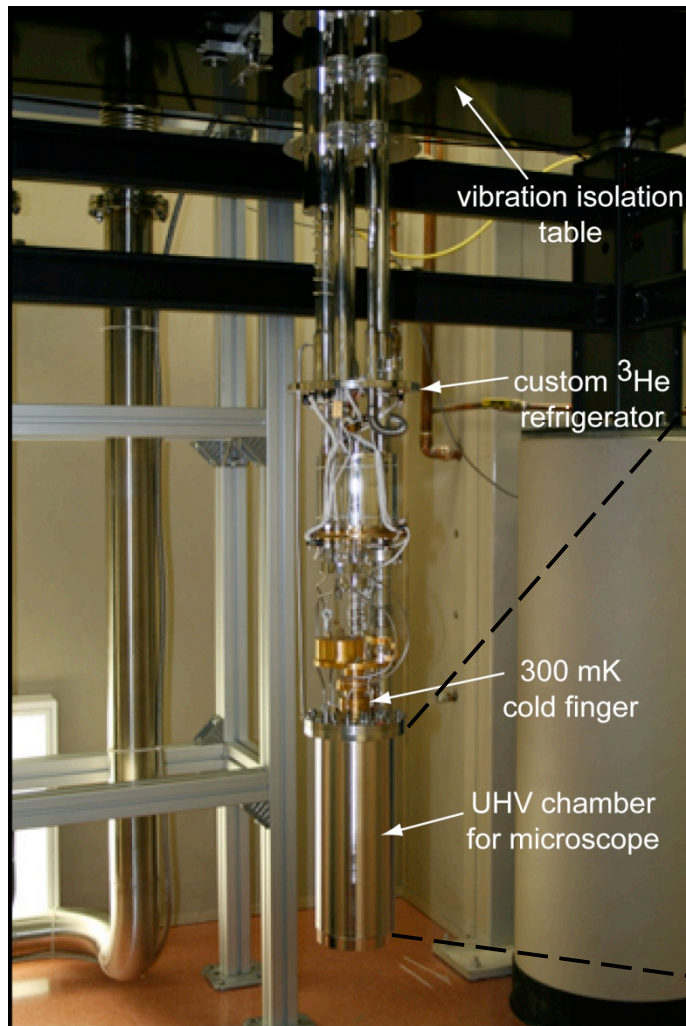
Cantilever Torque Magnetometry Measurements



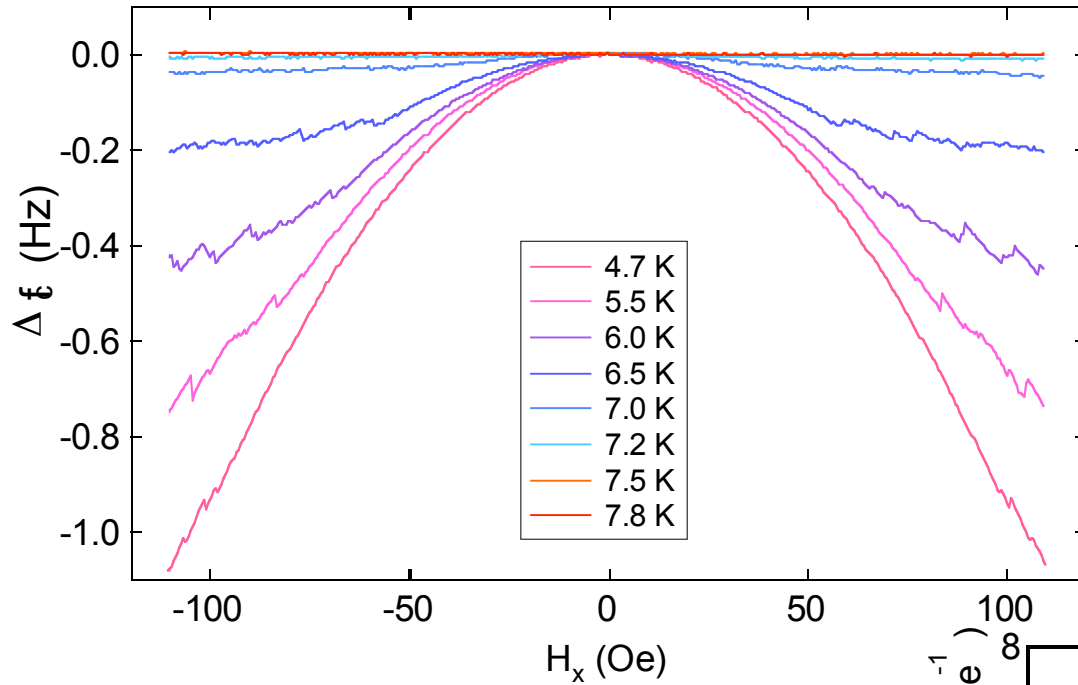
$$\left. \begin{aligned} f_0 &= 5.3 \text{ kHz} \\ k &= 3 \times 10^{-4} \text{ N/m} \\ Q &= 60,000 \\ l_{eff} &= 63 \text{ } \mu\text{m} \end{aligned} \right\} S_{\mu}^{1/2} = 3.3 \times 10^4 \mu_B / \sqrt{Hz} \text{ T}$$

$$T = 4.2 \text{ K}$$

300 mK Force Microscope

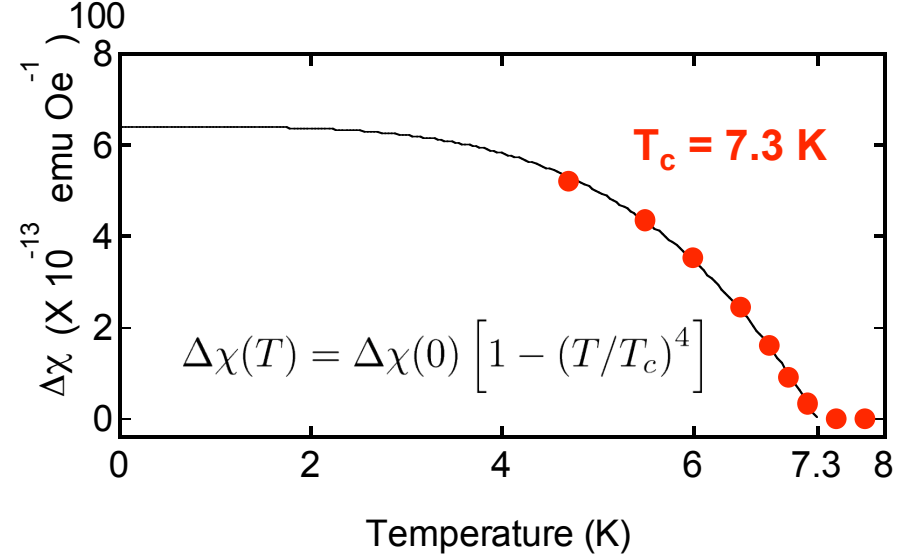
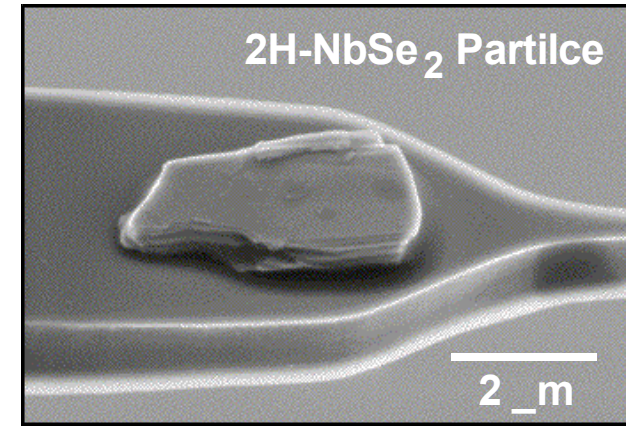


Micron-Size Superconducting NbSe₂ Particles

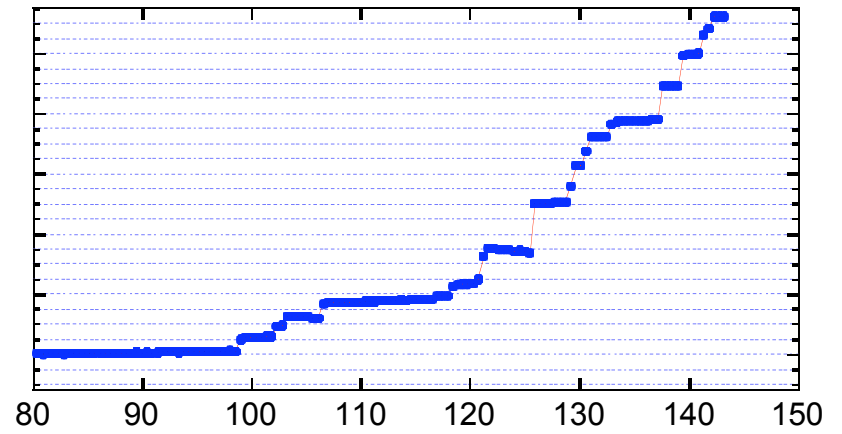
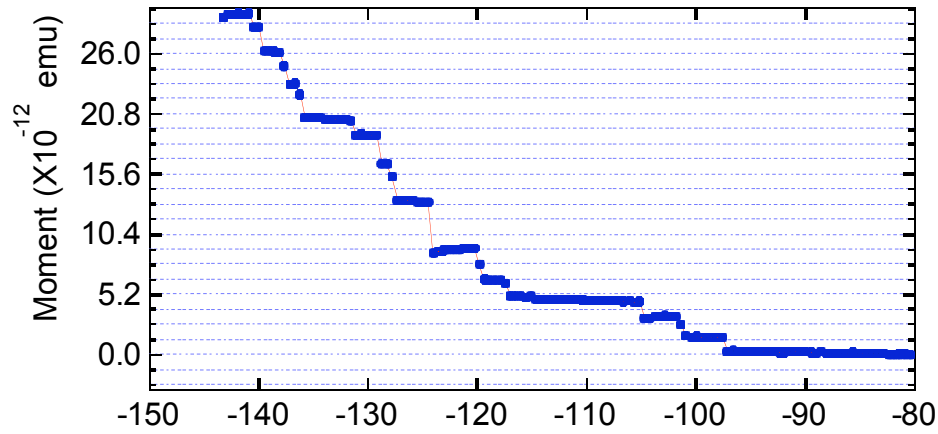
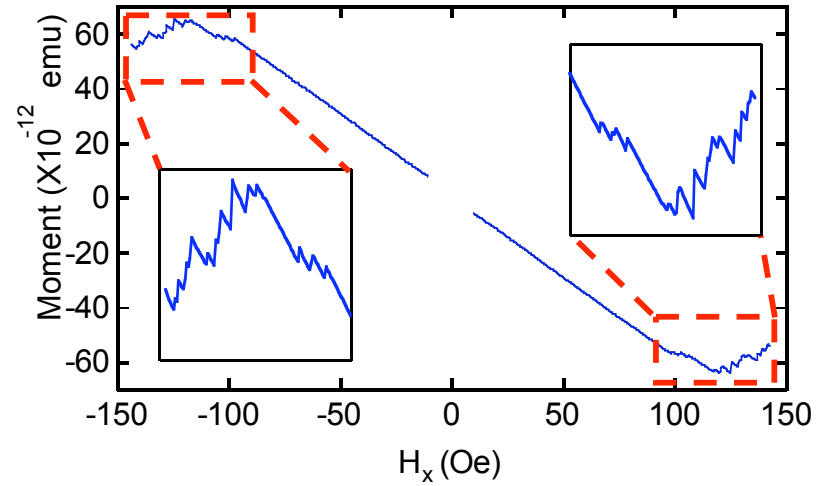
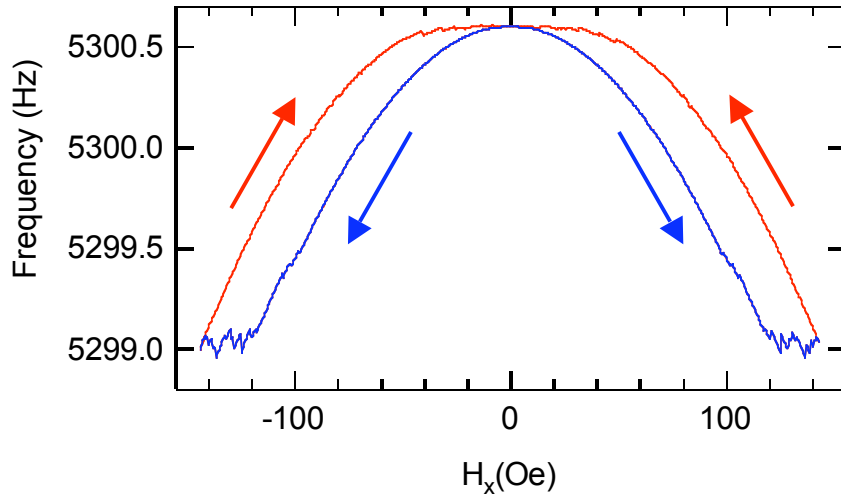


$$\lambda_c(0) = 230 \text{ nm}$$

$$\lambda_{ab}(0) = 69 \text{ nm}$$

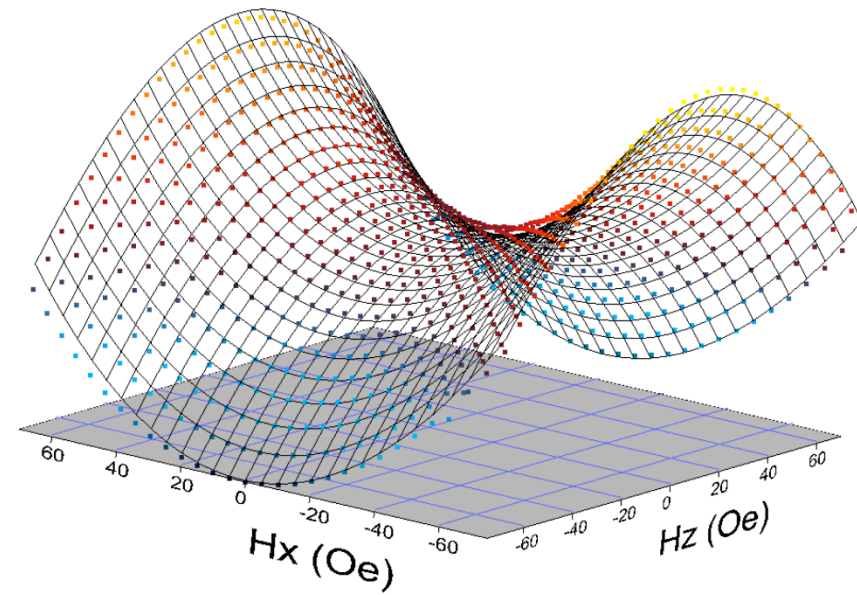
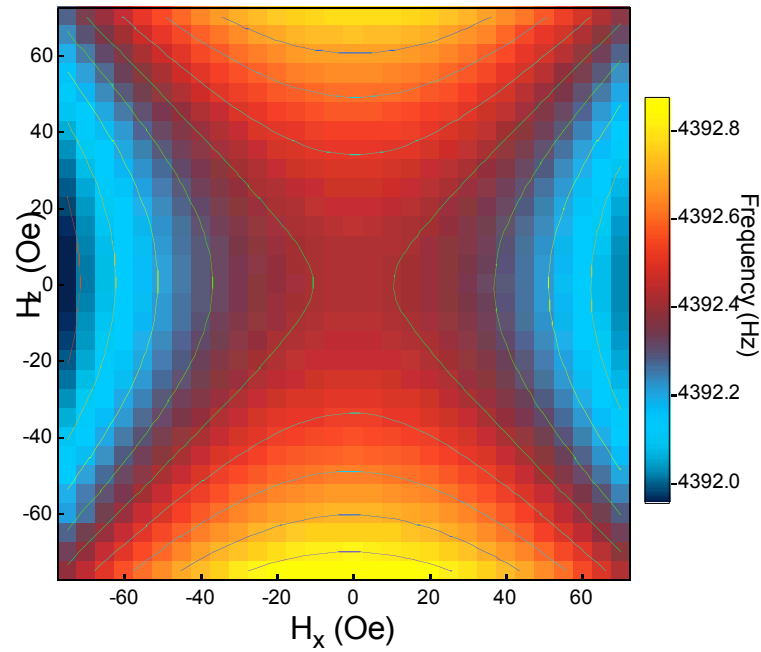


Vortex State in Mesoscopic NbSe₂ Samples



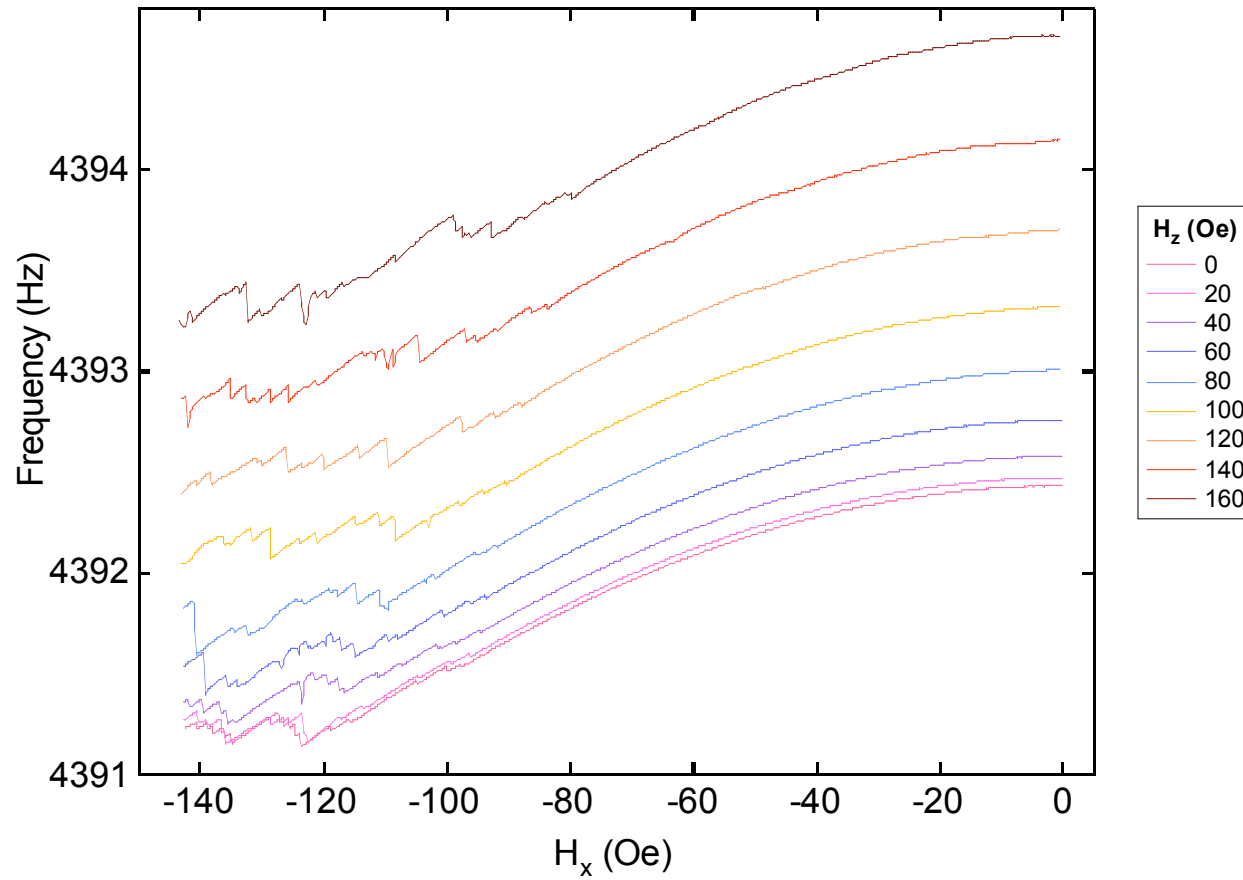
H_x (Oe)

Diamagnetic Susceptibility of NbSe₂

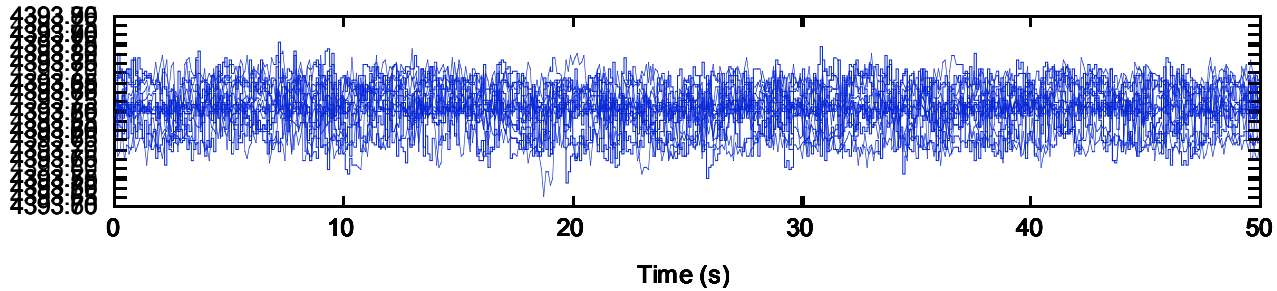
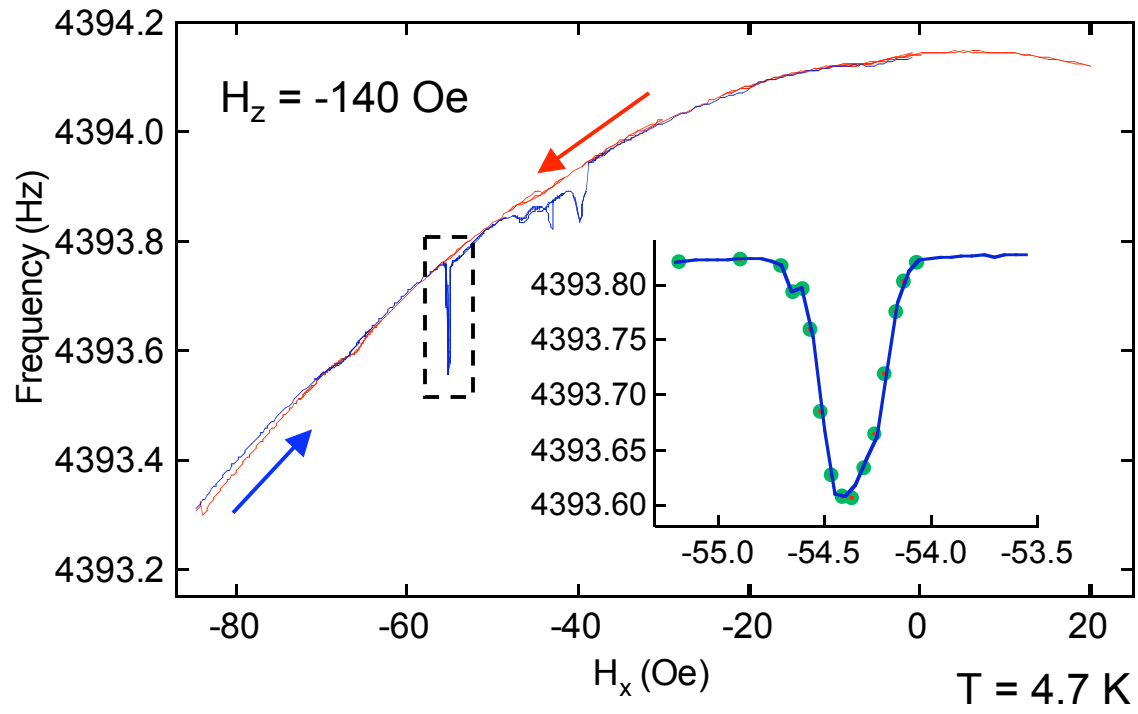


$$\Delta f = \frac{f_0}{2kl_{eff}^2} \Delta\chi (H_x^2 - H_z^2)$$

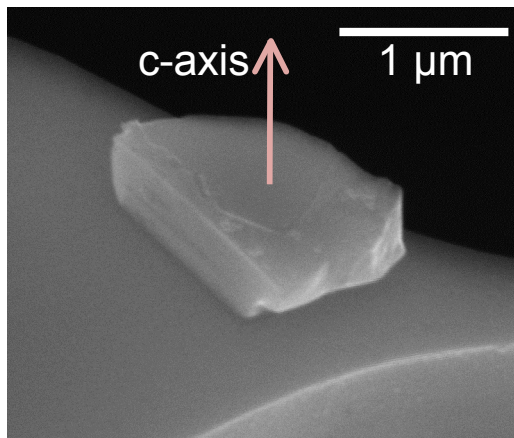
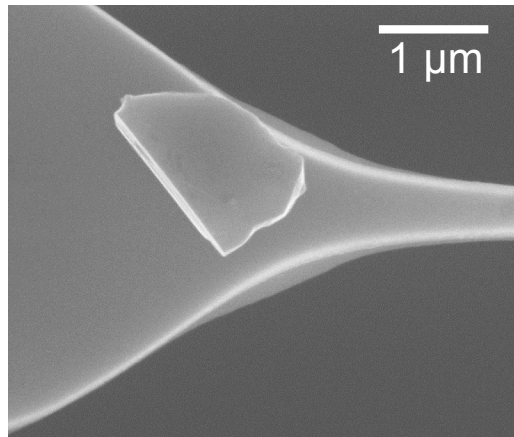
Response to ab-Plane Field



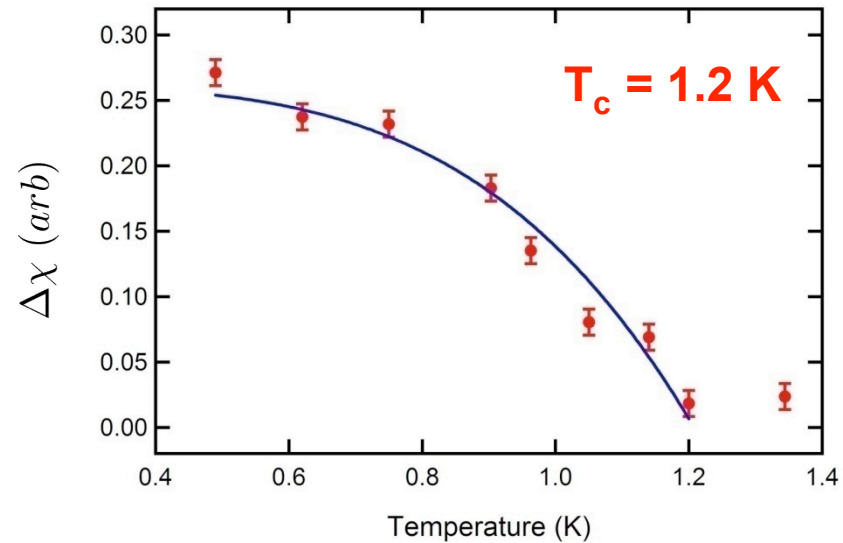
Switching Noise in Vortex Dynamics



Torque Magnetometry of Micron-Size Sr_2RuO_4 Particles



Samples grown by Y. Maeno



- Samples are cleaved from bulk crystals and glued to the cantilever with the c-axis normal to the cantilever face.

Parameter	<i>ab</i>	<i>c</i>
$\mu_0 H_{c2\parallel c}(0)$ (T)	0.075	
$\mu_0 H_{c2\parallel ab}(0)$ (T)	1.50	
$\mu_0 H_c(0)$ (T)	0.023	
$\xi(0)$ (Å)	660	33
$\lambda(0)$ (Å)	1520	3.0×10^4
$\kappa(0)$	2.3	46
$\gamma_s = \xi_{ab}(0)/\xi_c(0)$	20	

Zero-Field Magnetization Measurements

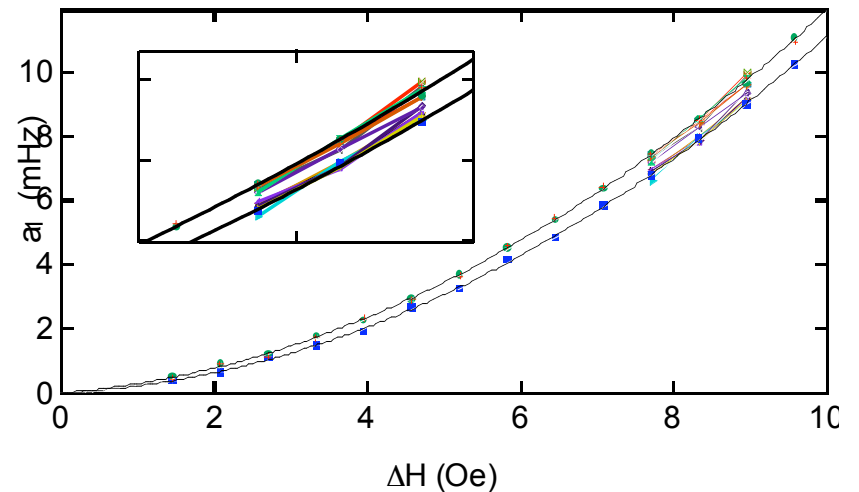
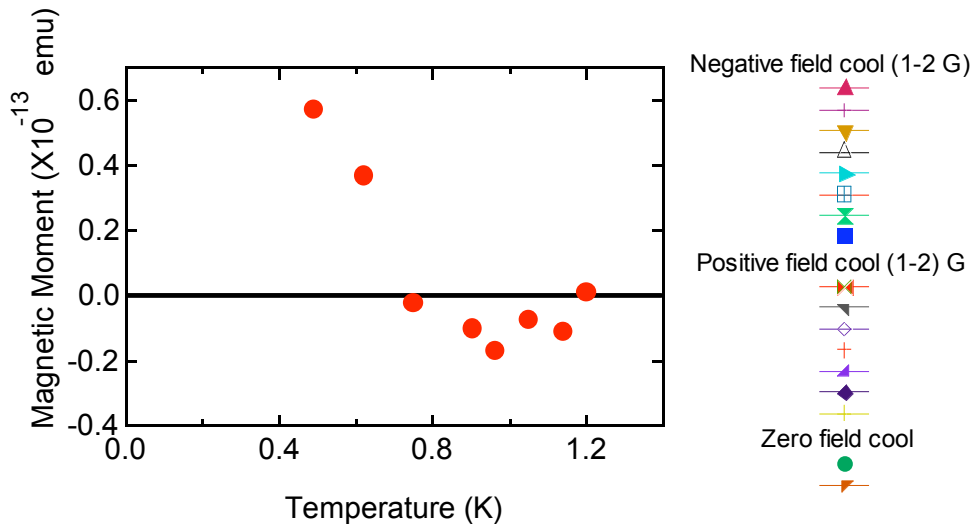
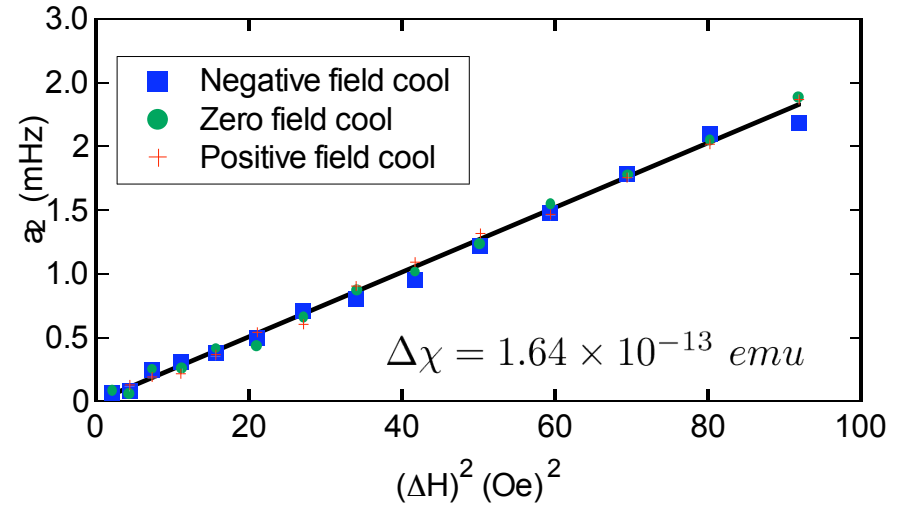
$$H_x = H_0 + \Delta H \cos(\omega_m t)$$

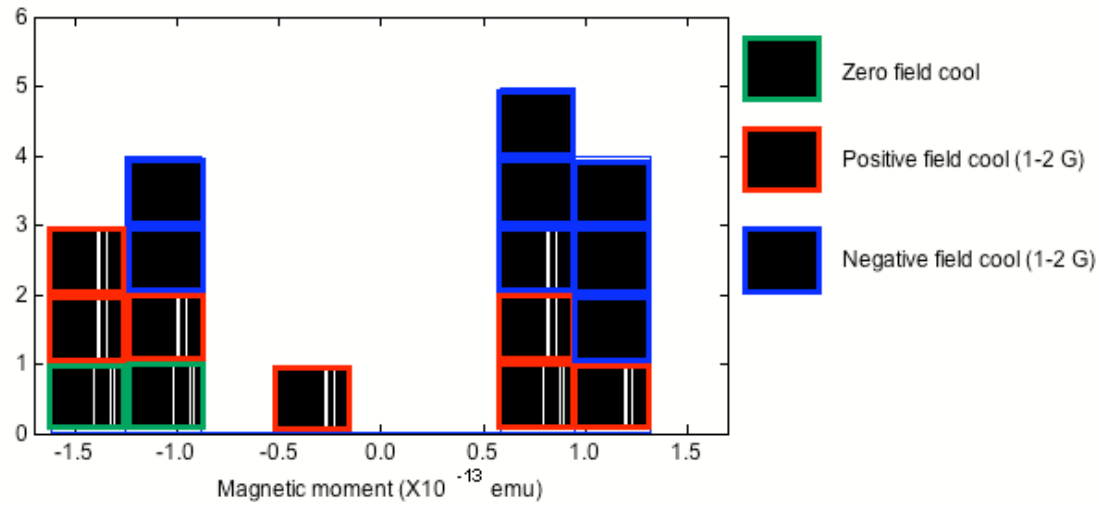
$$H_z = 0 \quad H_x^{min} = 1.25 \text{ Oe}$$

$$\Delta f = a_1 \cos(\omega_m t) + a_2 \cos(2\omega_m t) + const.$$

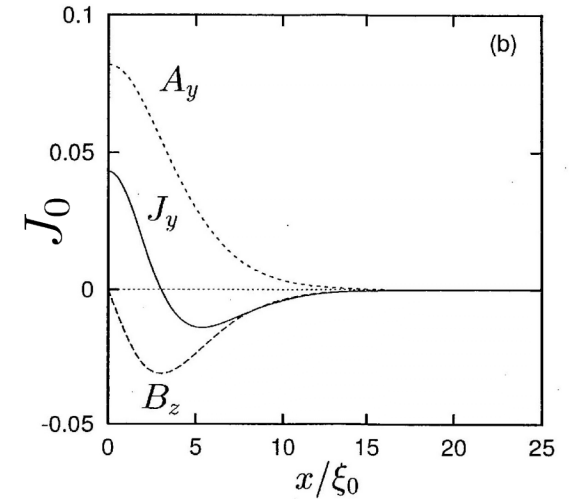
$$a_1 = \frac{f_0 \Delta H}{2kl_{eff}^2} (2H_0 \Delta\chi + \mu_x)$$

$$a_2 = \frac{f_0}{4kl_{eff}^2} \Delta\chi (\Delta H)^2$$



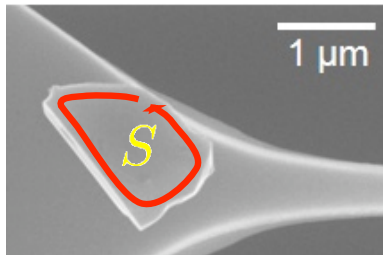


$$m_0 \approx (1.1 \pm 0.2) \times 10^{-13} \text{ emu}$$



M. Matsumoto & M. Sigrist
J. Phys. Soc. Jpn. **994** (1999)

Assume the particle is a single chiral domain



perimeter: $S \approx 4 \mu m$

thickness: $h = 440 \mu m$

$$m_0 \approx (0.5 - 1.0) \times 10^{-13} \text{ emu}$$

$$\xi_0 = 66 \text{ nm}$$

$$v_F = 5.5 \times 10^4 \text{ m s}^{-1}$$

$$N(E_F) = 4.36 \text{ states/eV cell}$$

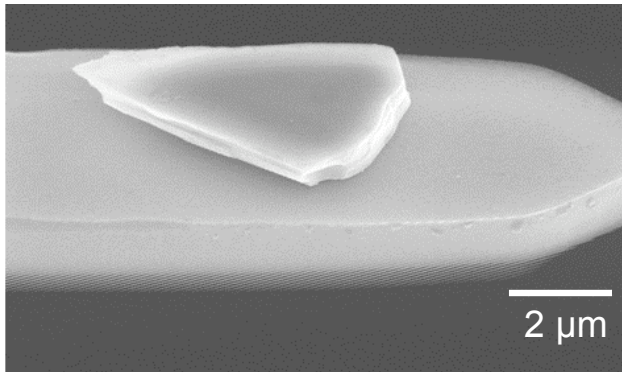
$$T_c = 1.5 \text{ K}$$

$$J_0 = ev_F N(E_F) k_B T_c$$

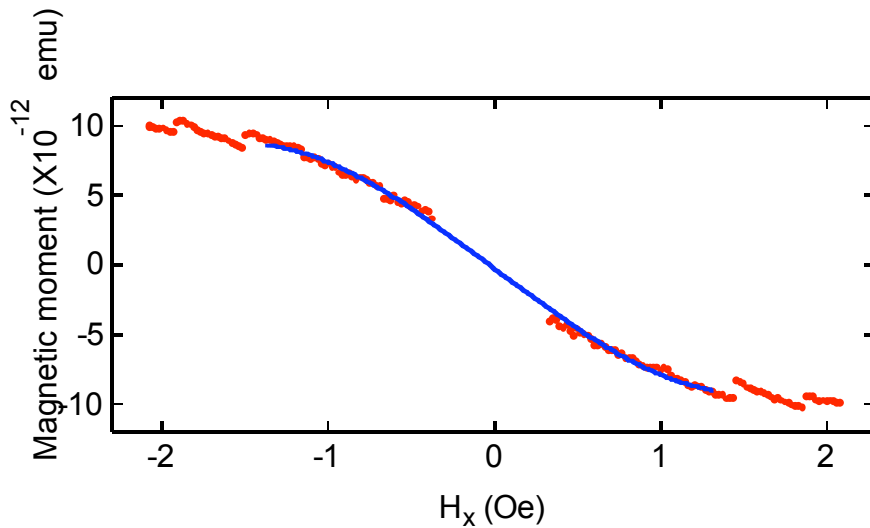
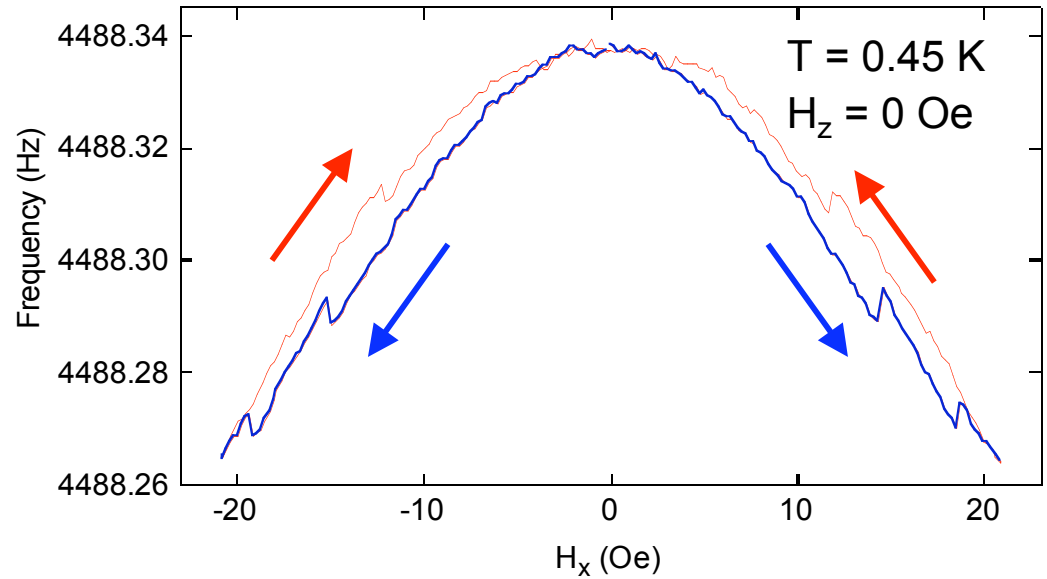
$$J_0 = 2.6 \times 10^{10} \text{ A m}^{-2}$$

Diamagnetic Susceptibility Measurements in Sr₂RuO₄

Sr₂RuO₄



Particle dimensions: 3 μm × 4 μm × 0.5 μm

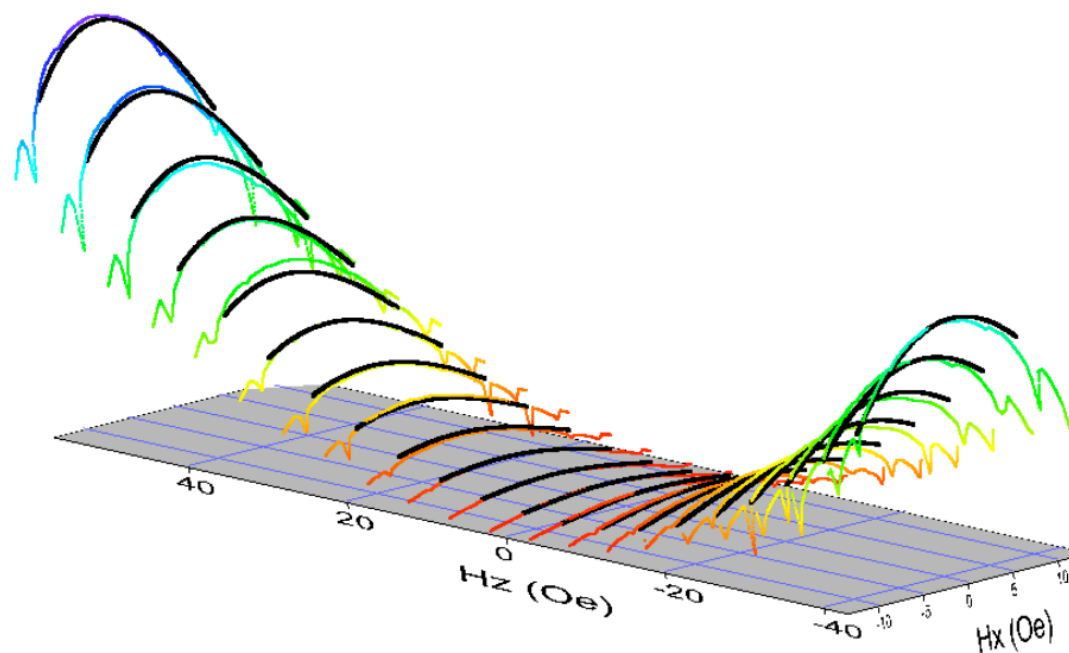
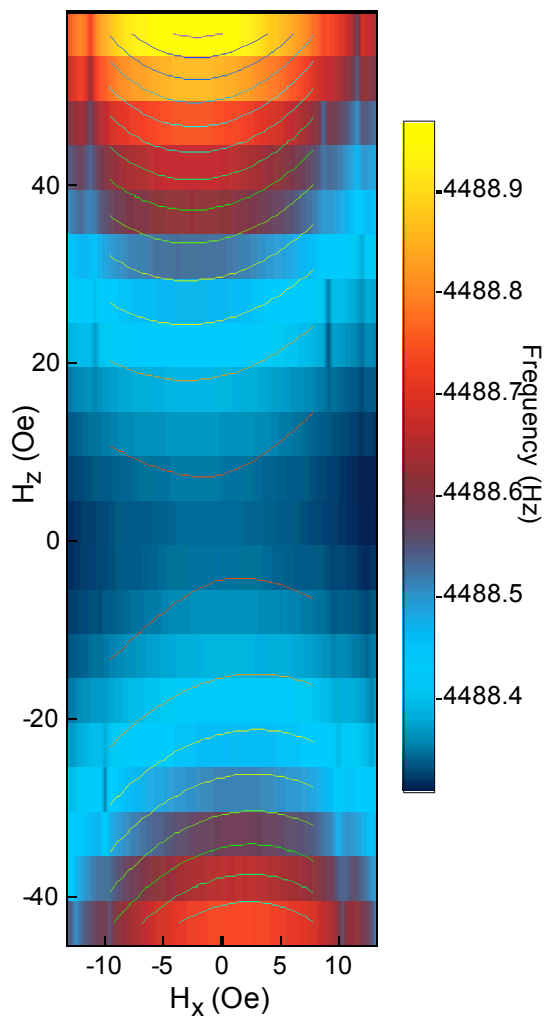


$$\Delta f \propto \Delta\chi_0 \left(1 - \frac{H_x^2}{H_0^2}\right) H_x^2$$

$$\Delta\chi_0 = -8.8 \times 10^{-13} \text{ emu}$$

$$H_0 = 26.7 \text{ Oe}$$





$$\Delta f \propto \left(1 - \frac{H_x^2}{H_0^2}\right) \Delta\chi_0 (H_x^2 - H_z^2)$$

$$\Delta\chi_0 = -5.5 \times 10^{-13} \text{ emu}$$

$$H_0 = 20.0 \text{ Oe}$$



Remarks

- Torque magnetometry measurements of mesoscopic samples is a promising technique for detection of edge currents
- Mesoscopic annular geometry might be useful in stabilizing fractional vortices

Questions

- Why do we not observe training effects in the zero-field magnetization ?
- What is the origin of the nonlinear diamagnetic susceptibility ?

