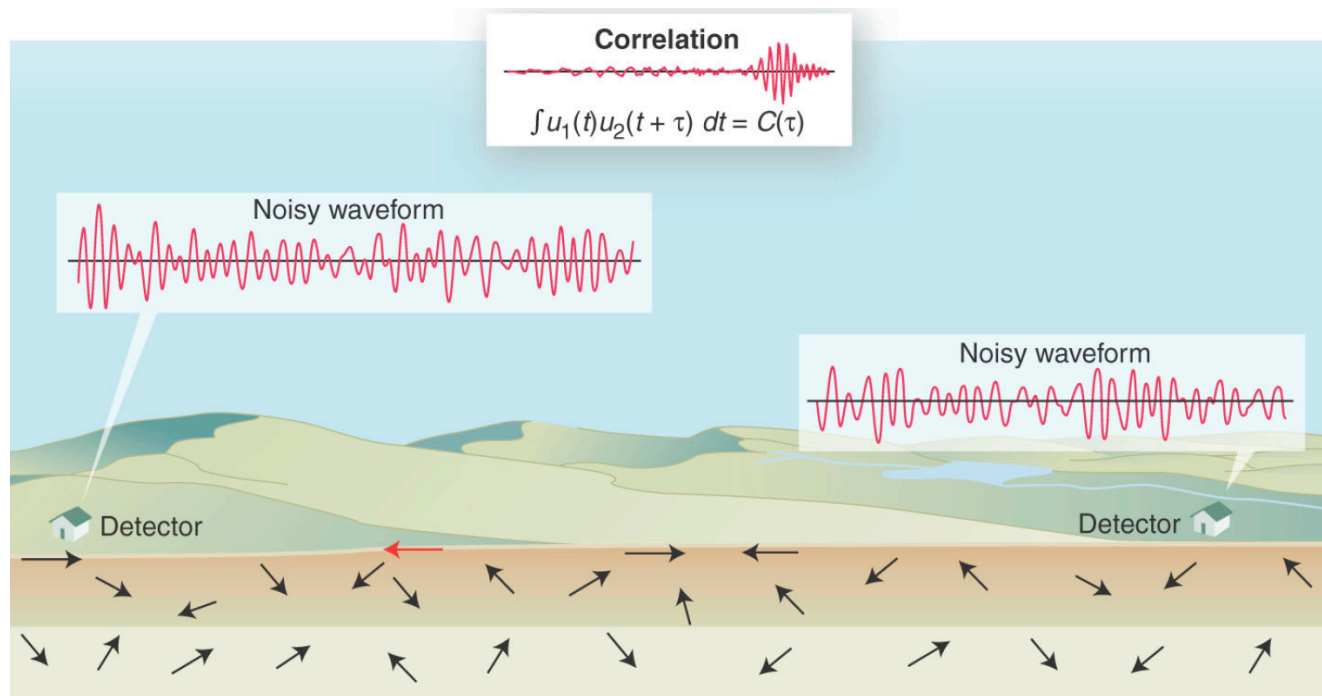


# Seismic Noise Correlations

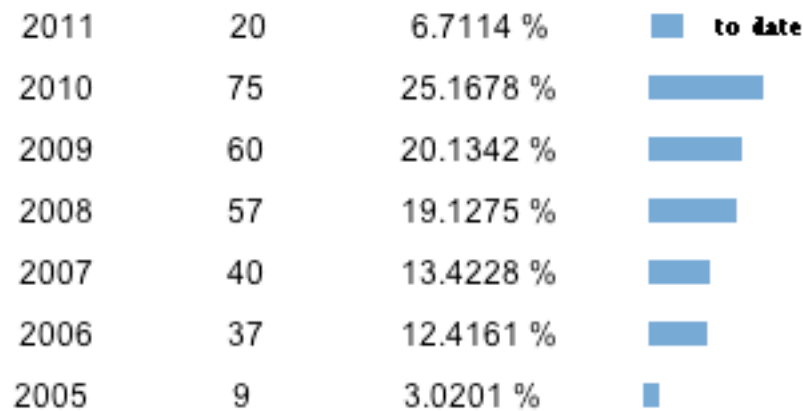
- *RL Weaver, U Illinois, Physics*



Over the last several years, Seismology

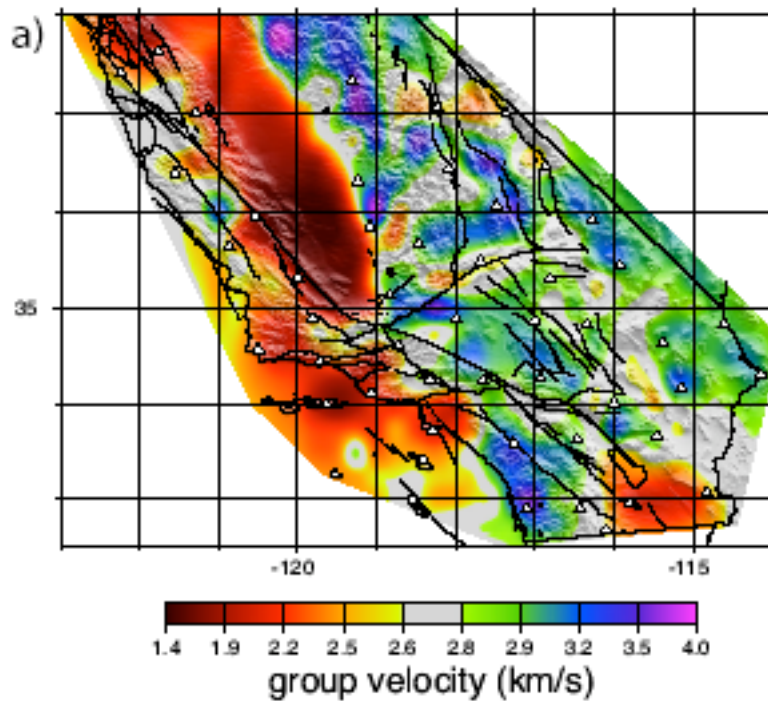
has focused growing attention on  
Ambient Seismic Noise  
and its Correlations.

Citation count on one of the seminal papers:



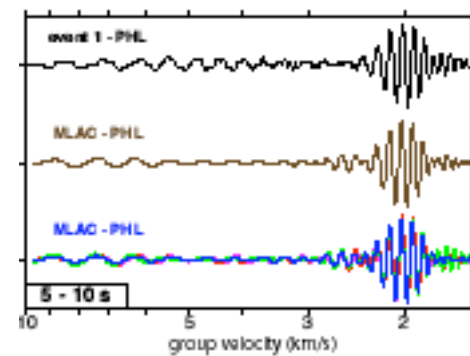
High-resolution surface-wave tomography from ambient seismic noise  
NM Shapiro *et al* *SCIENCE* 307 1615-1618 MAR 2005

The reason is (in part) due to the striking maps of seismic velocity that noise reveals . . .



A map of Surface-Wave Velocity in California

Obtained from correlating seismic noise



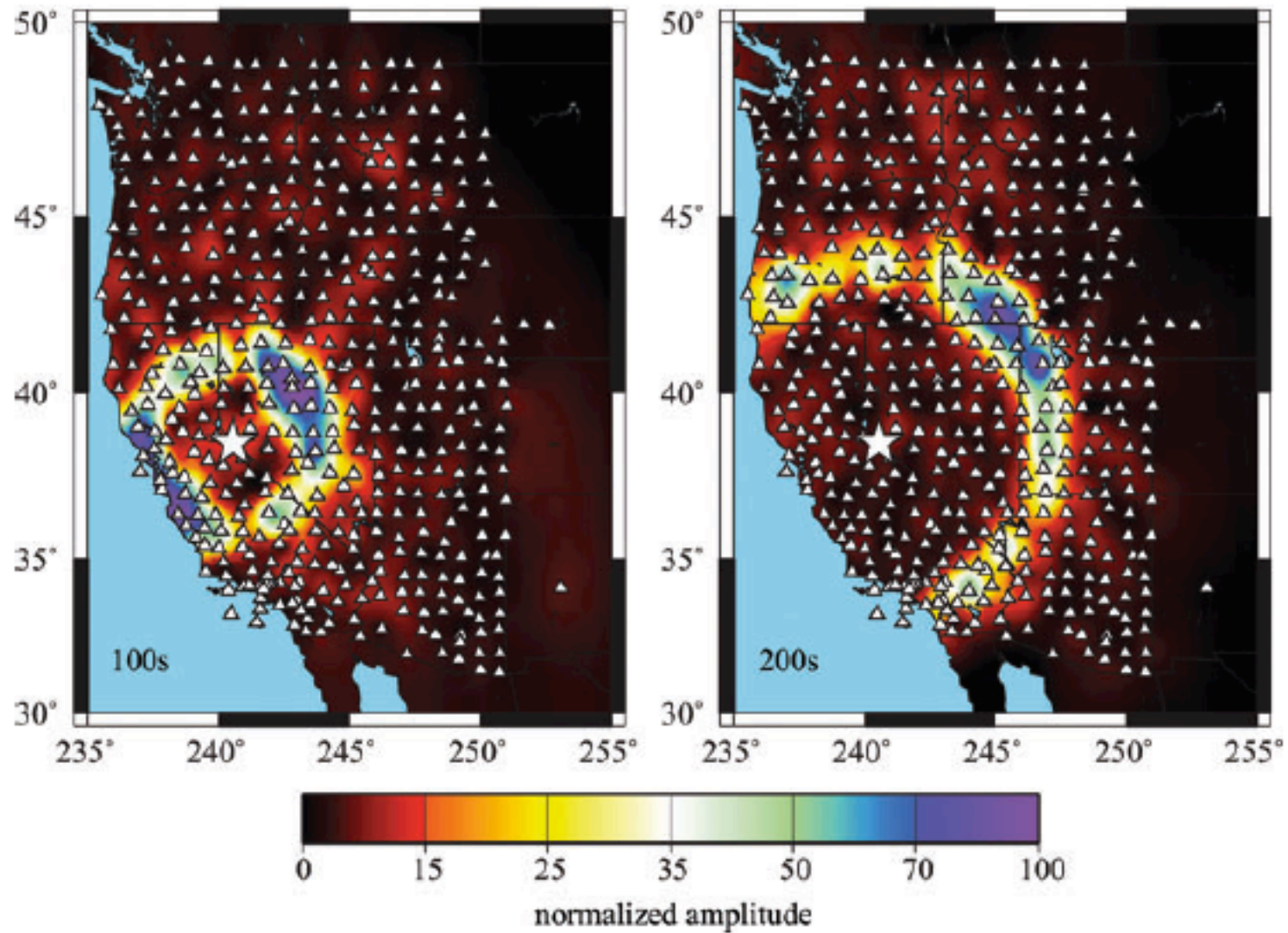
earthquake

1 year of correlations

4 one-month correlations

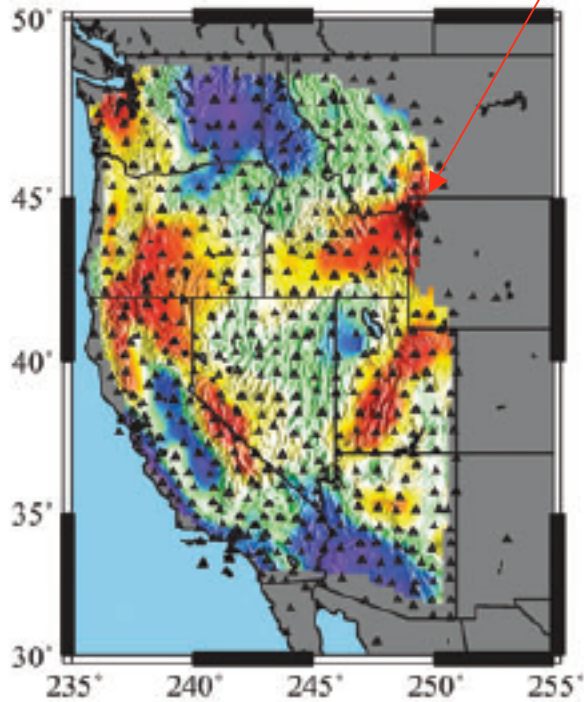
Frequencies  $\sim 0.02 < f < 1$  Hz;  $3\text{km} < \lambda < 150$  km

Lin and Ritzwoller and Snieder (2009) Geophys J Int  
3 years of data on a bigger array

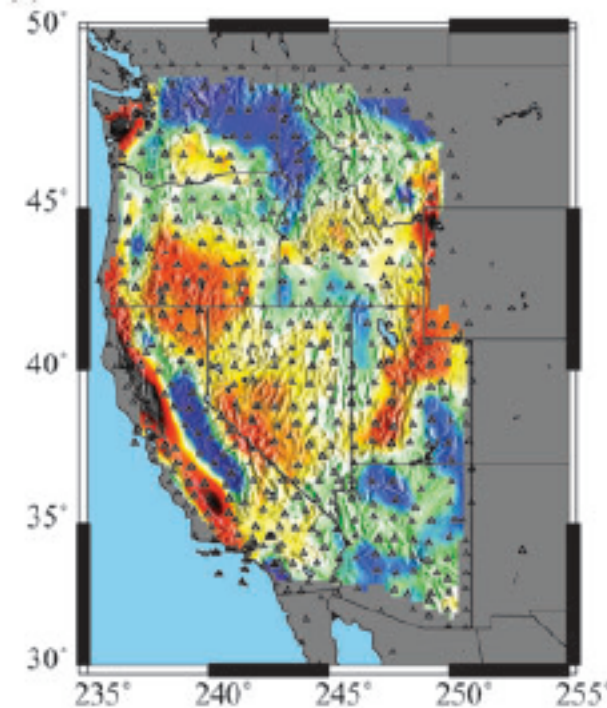
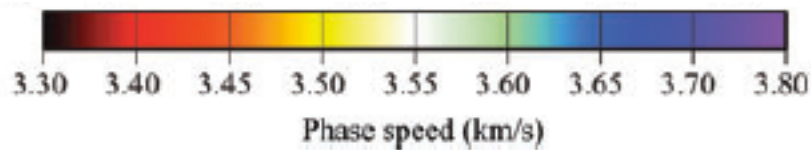


# Tomographically generated maps of wave speed

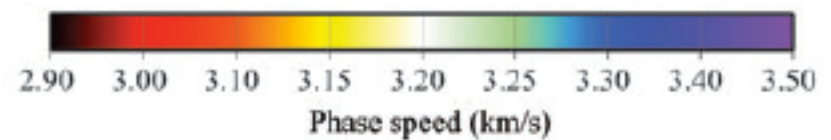
Hot spot in Yellowstone



24 sec

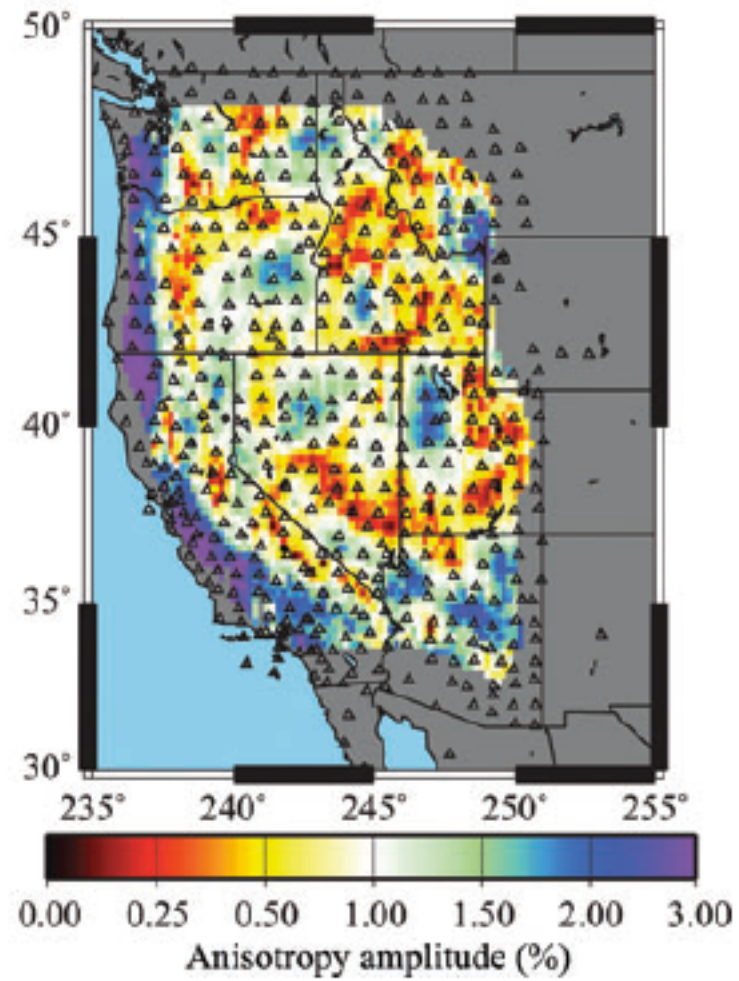


12 sec



Different properties at different frequencies  
i.e, different depths

They even resolve  $\sim 1\%$  anisotropies in wave speed



Main Assertion of Theory:

$$G(\vec{x}, \vec{y}; \tau) \sim \frac{\partial}{\partial \tau} \langle \psi(\vec{x}, t) \psi(\vec{y}, t + \tau) \rangle ?$$

Correlation of a diffuse field  $\psi(\vec{x}, t)$   
gives the Green's function  
( + sundry fine print and qualifications)

Where did this assertion come from?

Why should we believe it?

How *much* should we believe it?

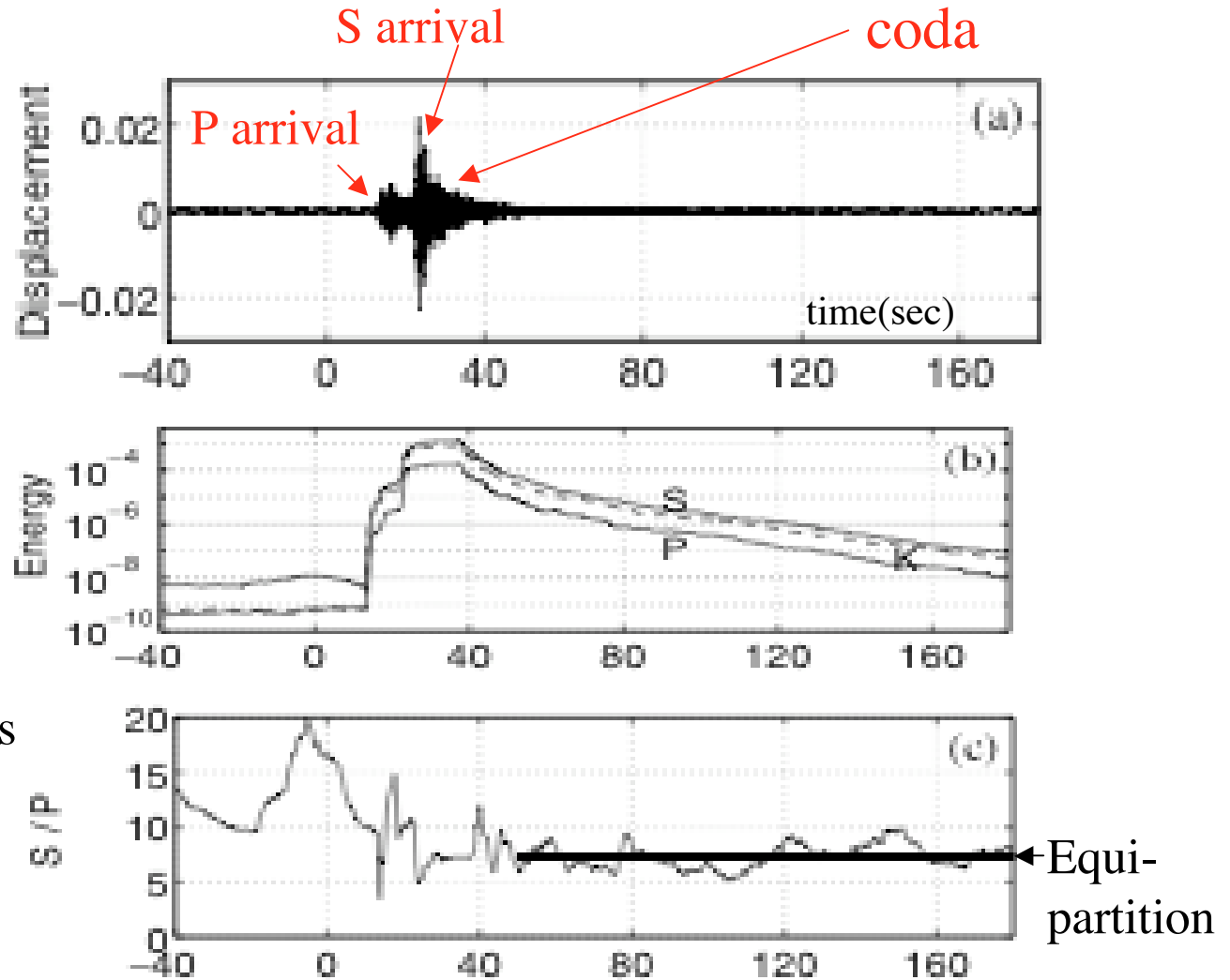
## History of the approach . . .

Conversations with a seismologist, at a 1999 workshop, about the seismic coda - which appeared to be equipartitioned

An earthquake record

Ray arrivals are followed by low amplitude noise, or "*coda*"

The coda appears to achieve a steady state ratio of its energy contents  
For example, its shear-to-dilatational energies: S/P





## History of the approach continued. . .

I then pointed out

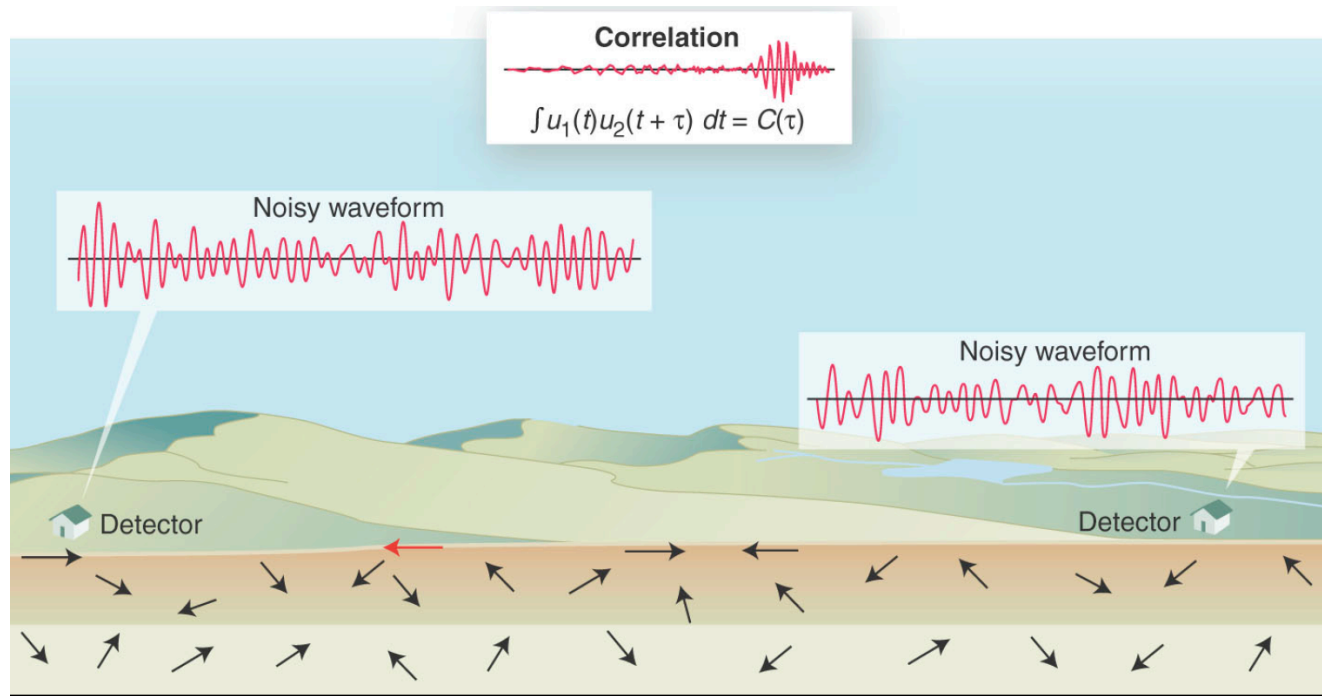
"if a wave field (e.g. seismic coda) is multiply scattered to the point of being equipartitioned, the field's correlations should be Green's function,

And we could recover lots of information without using a controlled source"

Geophysicist: "useful, if true"

Physicist: "Nonsense, can't possibly be true"

## Hand-waving plausibility argument . . .



that there could well be a signature, an "arrival,"  
at the correct travel time

- due to those few rays that happen to be going the right way

But G exactly?

And where's the proof?

And won't other ray directions obscure the effect?

## Standard Proofs . . .

- For a thermally diffuse field
  - modal picture
  - fluctuation-dissipation theorem
- For a conventional acoustic diffuse field
  - modal picture ( sensible only for closed systems )
  - plane wave picture (sensible only for homogeneous systems)
- Systems with uniformly distributed incoherent sources everywhere
- Heterogenous loss-free region without sources,  
but *insonified* by an external diffuse field

## But what about imperfectly diffuse fields?

- There is an asymptotic ~validity to assertion

The simplest proof involves a common definition of a fully diffuse field, from room acoustics or physics of thermal phonons:

in terms of the normal mode expansion for the field in a finite body

$$\phi(x, t) = \Re \sum_{n=1}^{\infty} a_n u_n(x) \exp\{i\omega_n t\}$$

$$\langle \mathbf{a}_n \mathbf{a}_m^* \rangle = \delta_{nm} F(\omega_n) / \omega_n^2 \quad \text{"equipartition"}$$

*n.b: this follows from maximum entropy  
where  $F \sim$  energy per mode ( $k_B T$ )*

$$C \equiv \langle \phi(x, t) \phi(y, t+\tau) \rangle = \frac{1}{2} \Re \sum_{n=1}^{\infty} F(\omega_n) u_n(x) u_n(y) \exp\{-i\omega_n \tau\} / \omega_n^2$$

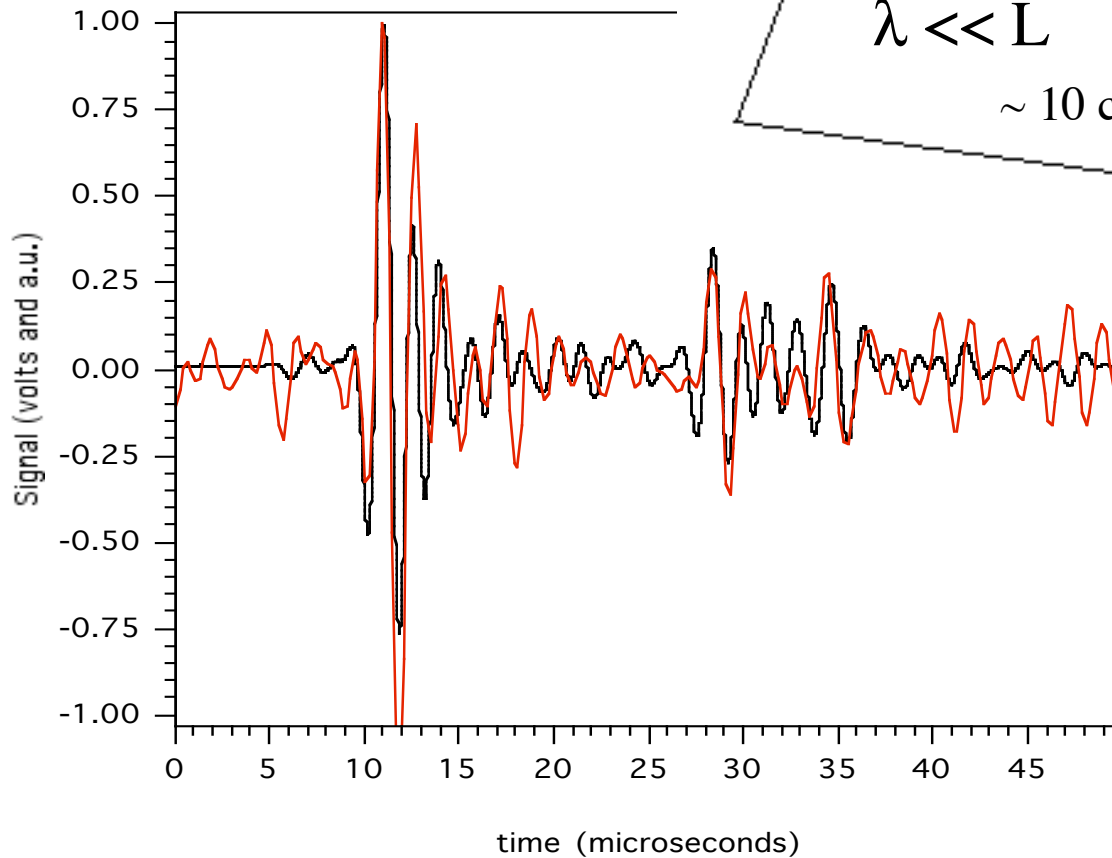
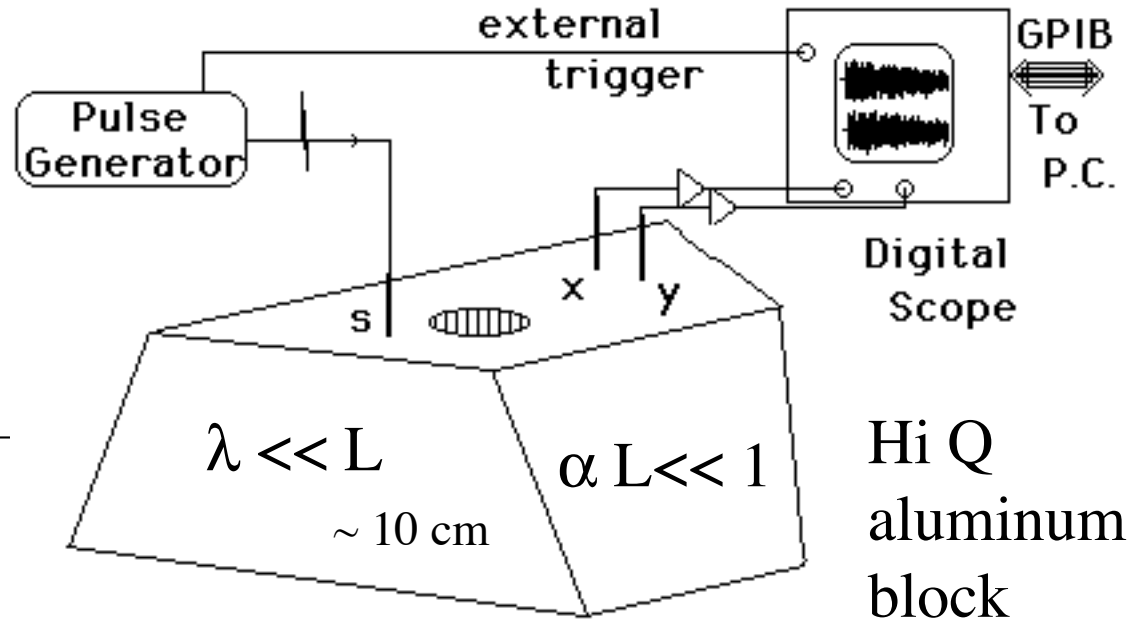
Compare with  $G \dots$

$$G_{xy}(\tau) = \sum_{n=1}^{\infty} u_n(x) u_n(y) \frac{\sin \omega_n \tau}{\omega_n} \quad [\text{for } \tau > 0, 0 \text{ otherwise}]$$

So,  $\partial C / \partial \tau = G - G^{\text{time reversed}}$ , i.e.,  $G - G^*$  or  $\text{Im } G$  if  $F$  is constant

Verification?

e.g. . . .



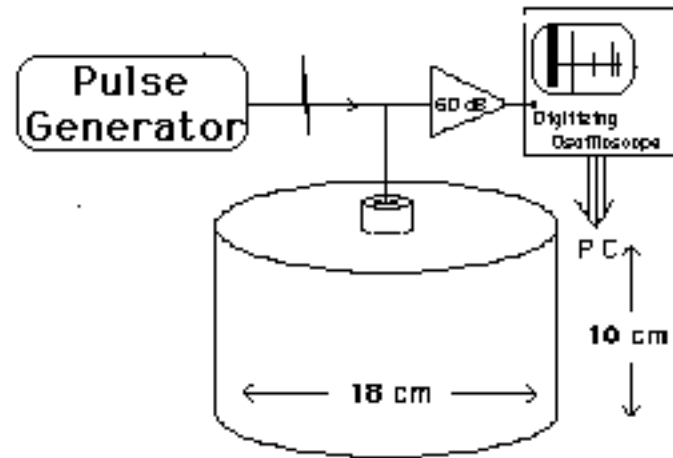
— direct pitch/catch signal

— correlation

*J Acoust Soc Am.* **110**,  
(2001)

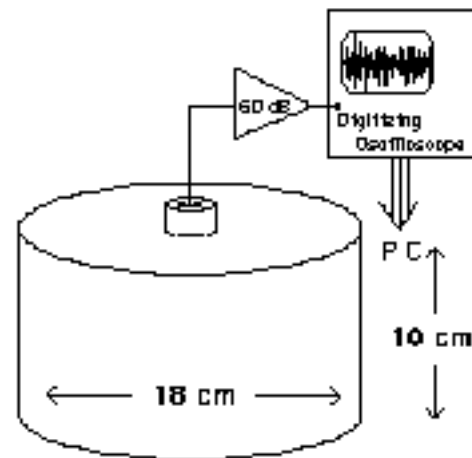
Comparison of a  
Direct Pulse-Echo  
Signal,

(conventional ultrasonics)



and

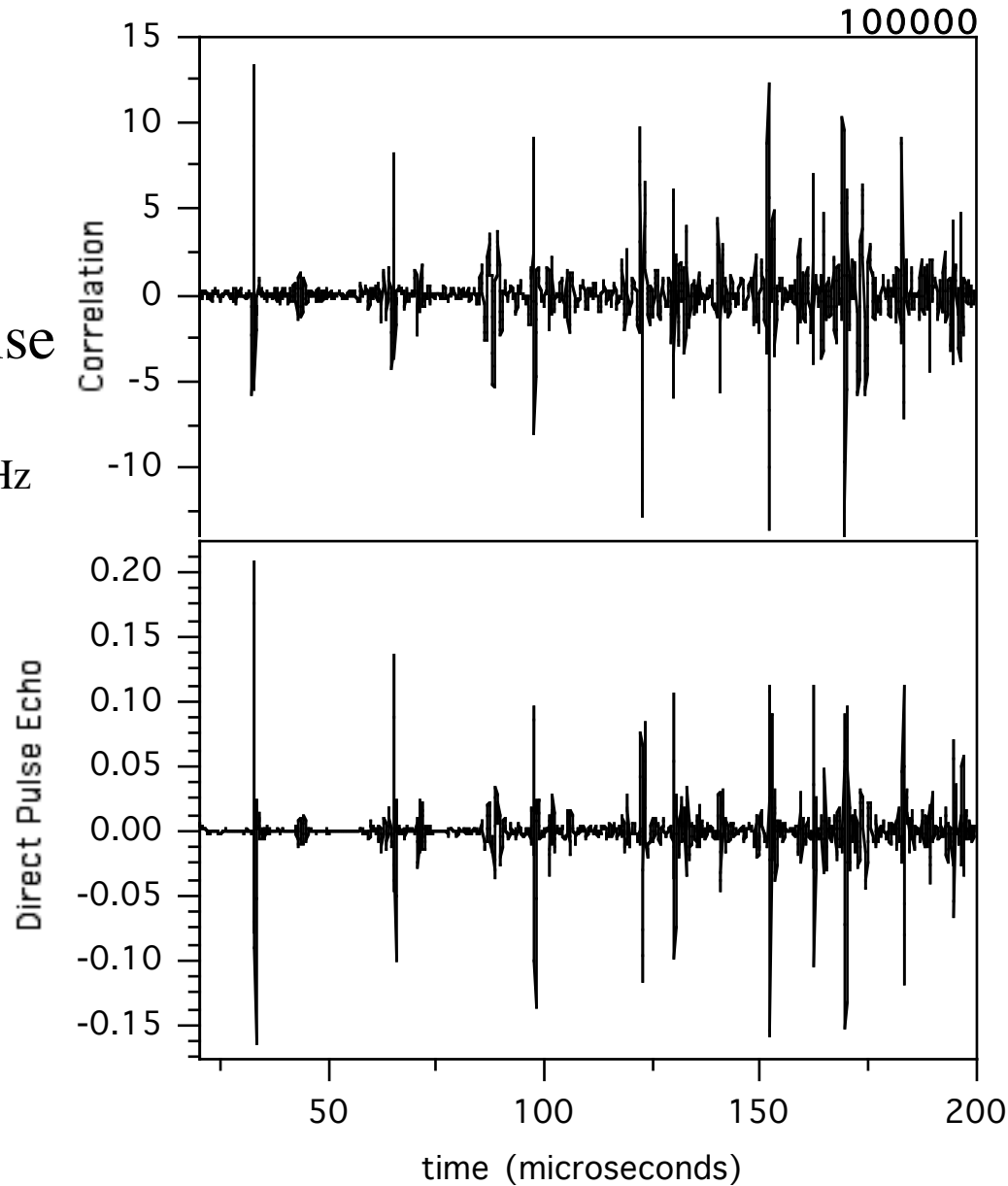
Thermal Noise  
Correlation



*Phys Rev Lett* **87** 134301 (2001)

Correlation  
Of  
Thermal noise

rms  $u \sim 3 \text{ fm}/\sqrt{\text{MHz}}$



After  
Capturing  
320 seconds  
Of data

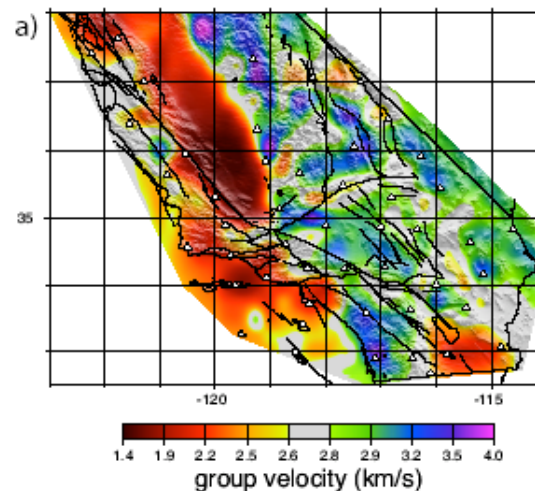
(and taking 2.5  
hours to do so)

But proofs that require full diffusivity  
and/or finite bodies and closed acoustic systems,  
May not be relevant for practice.

Ambient seismic noise(\*), for example, is  
NOT fully diffuse

It has preferred directions (sources in ocean storms)

Nevertheless, these  
maps are impressive



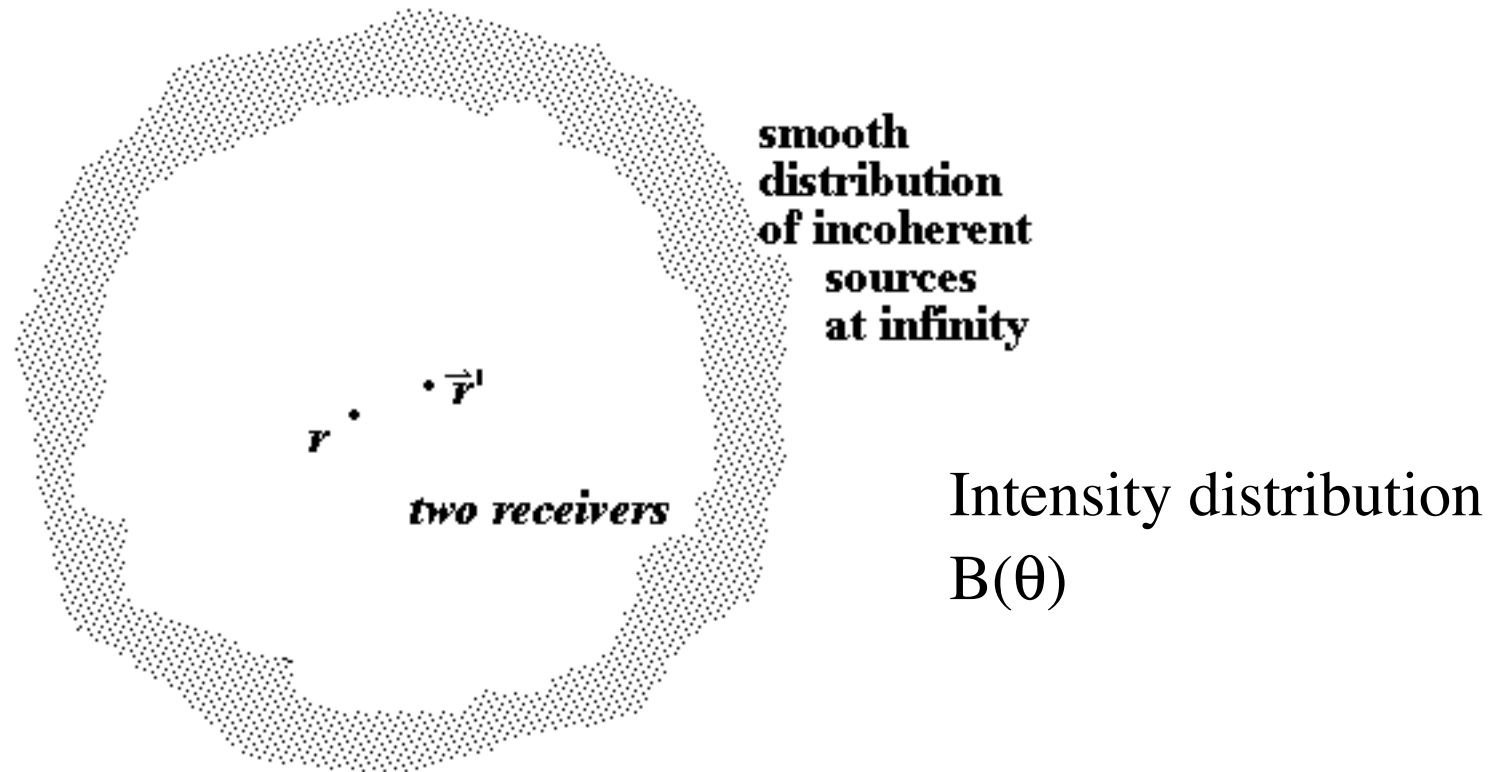
Why does it work?

\*Late *coda* appears  
fully diffuse, but  
there isn't enough of it.



What if an incident field does not have isotropic intensity?  
What if it is not equipartitioned?

Consider a homogeneous medium with  
incoherent sources at infinity



The field in the vicinity of the origin is a superposition of plane waves

$$\psi(\vec{r}, t) = \int A(\theta) \exp(-ik\hat{\theta} \cdot \vec{r} + i\omega t) d\theta \quad (2-d)$$

with  $\langle A \rangle = 0$ ;  $\langle A(\theta)A^*(\theta') \rangle = B(\theta)\delta(\theta - \theta')$

i.e, an incident plane wave intensity  $B(\theta)$



Which implies that the field-field correlation is

$$\langle \psi(\vec{r}, t)\psi(\vec{r}', t') \rangle = \int B(\theta) \exp(-i\omega\hat{\theta} \cdot (\vec{r} - \vec{r}') / c + i\omega(t - t')) d\theta$$

exact

$$C = \langle \psi(\vec{r}, t) \psi(\vec{r}', t') \rangle = \int B(\theta) \exp(-i\omega \hat{\theta} \cdot (\vec{r} - \vec{r}') / c + i\omega(t - t')) d\theta$$

wavelet  $S(t)$  related to power spectrum of noise

$$= \frac{-1}{4\pi} \sqrt{\frac{2\pi}{\omega x}} \int_0^{+\infty} d\omega i \exp(i\omega(t - x/c)) \tilde{S}(\omega - \omega_0) \times$$

$$\{ B(0) e^{i\pi/4} + B''(0) \frac{1}{2\omega x} e^{3i\pi/4} - B(0) \frac{i}{8\omega x} e^{5i\pi/4} \dots \} + c.c.$$

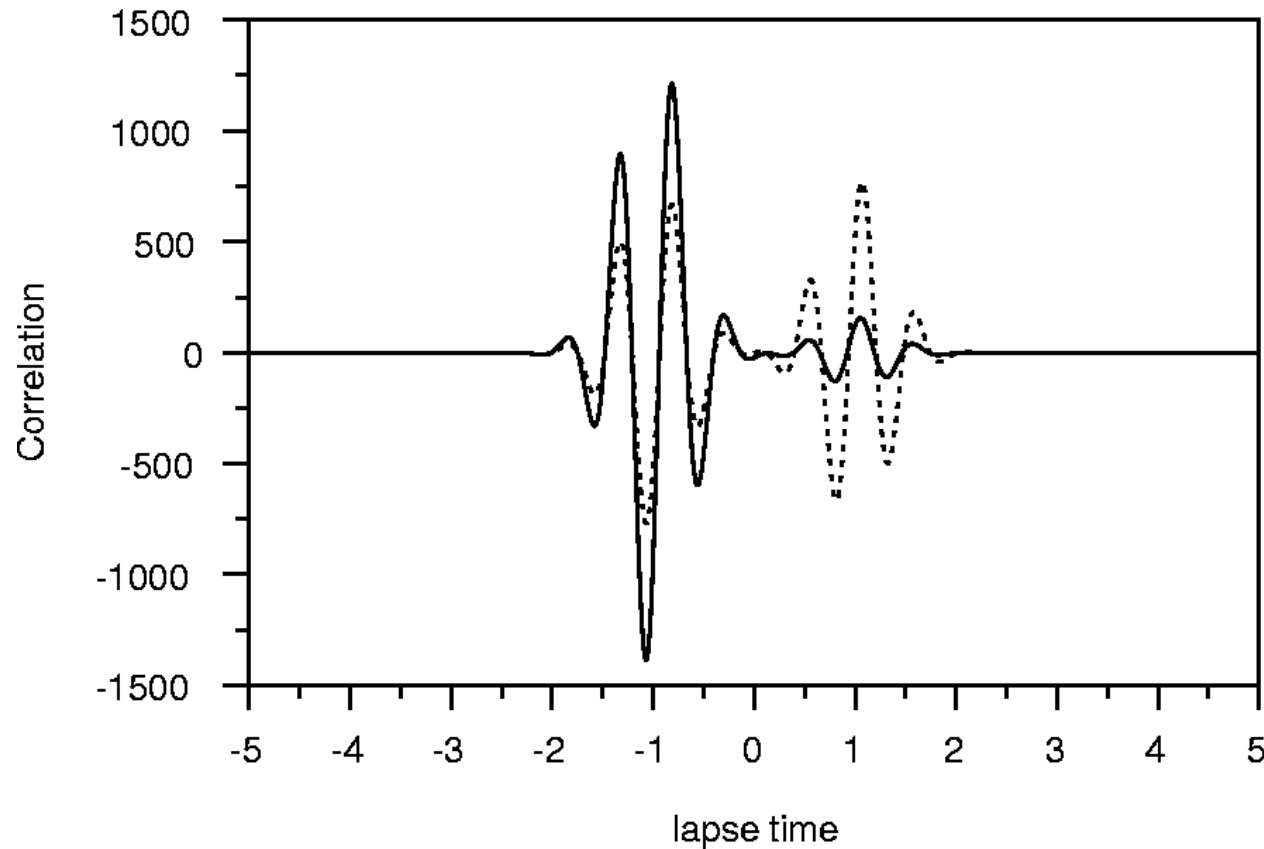
Leading term  $\nearrow$   $\nwarrow$  first correction

Permits us to show that the apparent arrival time is delayed relative to  $|r-r'|/c$  by a fractional amount  $B''(0)/2k^2|r-r'|^2B(0)$

$\Rightarrow$  The effect of non-isotropic  $B$  on arrival time is small in practice  
 $\Rightarrow$  Hence the high quality of the maps of seismic velocity even though the ambient seismic noise is not equipartitioned

Comparison of Correlation waveform (solid line)  
and time-symmetrized G (dashed line)

For case of non-trivial ponderosity  $B(\theta) = 1 - 0.8 \cos \theta$



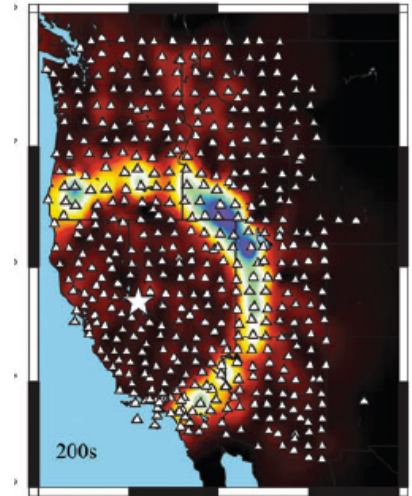
Froment et al  
2009

- Note:
- a) assertion fails,  $G \neq dC/d\tau$
  - b) large differences in positive and negative time amplitudes
  - c) there are *tiny* shifts of apparent arrival time, as predicted

In sum, the method works well for arrival times, hence the good maps

The method is well suited to seismology because

Controlled sources are highly inconvenient,  
(earthquakes and nuclear explosions)



Recent advent of large arrays of long-period seismic stations  
and world-side access to their time records

For many years seismologists would record seismic time-records,  
ignore the noise, and examine the earthquakes

Now they throw out the earthquakes and keep the noise.

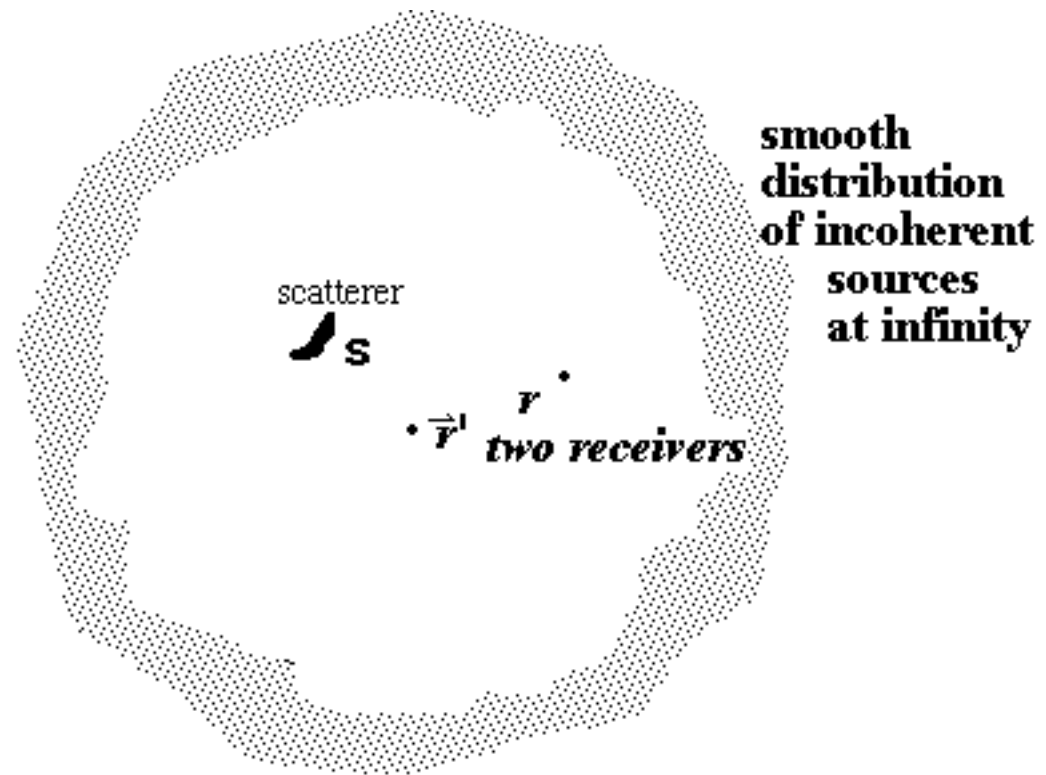
Other consequences of imperfectly partitioned ambient noise:

Spurious features in the correlations due to scatterers

Amplitude information is hard to interpret

Spurious arrivals..

Intensity distribution  
 $B(\theta)$



Correlations in the presence of a scatterer will show

a direct arrival at  $\tau = |\mathbf{r} - \mathbf{r}'|/c$

an indirect arrival at  $\tau = |\mathbf{r} - \mathbf{s}|/c + |\mathbf{r}' - \mathbf{s}|/c$

*and*

a spurious arrival at  $\tau = |\mathbf{r} - \mathbf{s}|/c - |\mathbf{r}' - \mathbf{s}|/c$

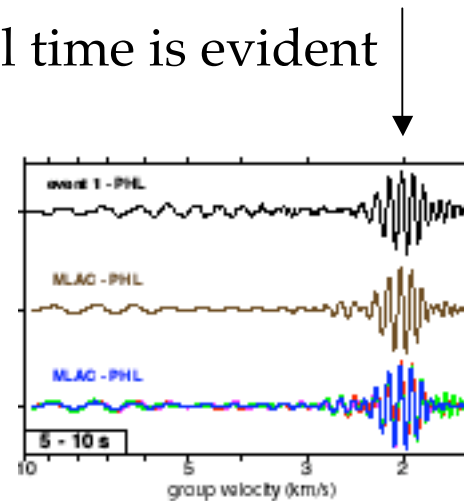
} parts of  $G$

Disappears if field is  
equipartitioned

## Amplitude information?

The technique has been used very successfully in seismology to recover seismic velocities, with high spatial resolution. But . . .

Arrival time is evident



Arrival amplitude?  
Is this meaningful?

If we really had  $G$ , we'd be able to infer attenuation also.

Issues include the unknown field intensity in the direction between the detectors



Ray amplitudes  $X$  depend on  
attenuation  $\alpha$   
"on-strike" intensity  $B$

$$X_{i \rightarrow j} = B_i(\hat{n}_{i \rightarrow j}) \sqrt{2\pi / \omega_o |\vec{r}_i - \vec{r}_j|} \exp\left(-\int_{\vec{r}_i}^{\vec{r}_j} \alpha(\vec{r}) d\ell\right)$$

If field is not equipartitioned, then  $B$  varies. But how?

Noise intensity  $B$  varies in space like an RTE?

$$\hat{n} \cdot \vec{\nabla} B(\vec{r}, \hat{n}) + 2\alpha(\vec{r}) B(\vec{r}, \hat{n}) = \underbrace{P(\vec{r}, \hat{n})}_{\text{sources}} + \underbrace{\oint B(\vec{r}, \hat{n}') p(\vec{r}, \hat{n}, \hat{n}') d\hat{n}'}_{\text{Scattering into direction n}}$$

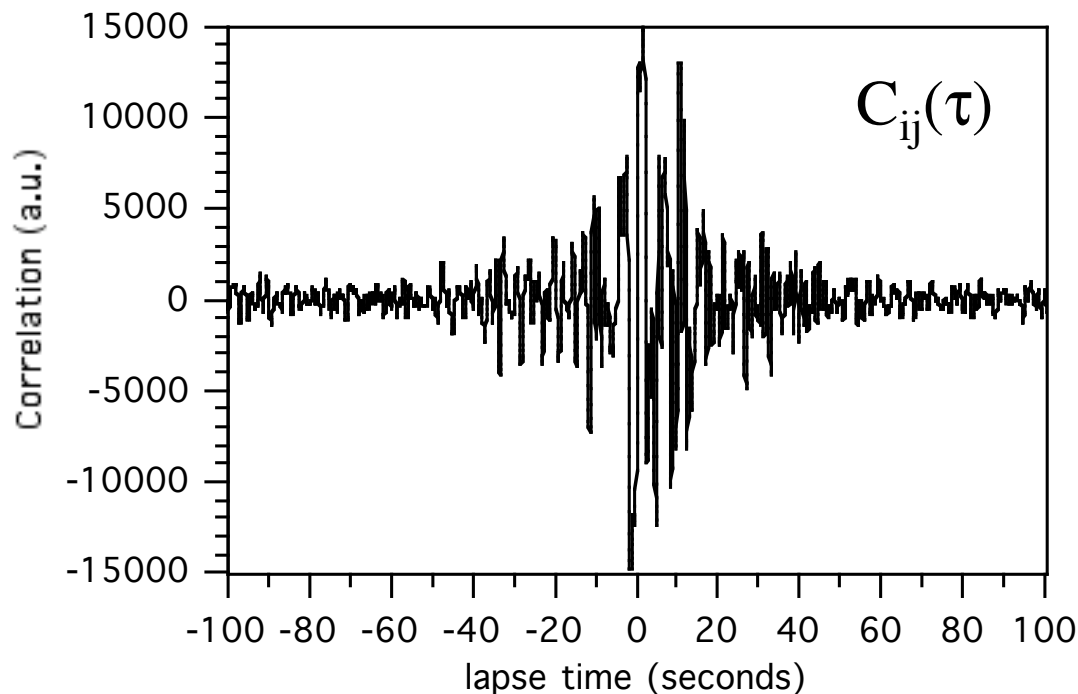
## Another application

Detecting changes in a medium . . .

Brenguier et al *Science*: **321**. 1478 - 1481(2008)

correlated ocean-generated seismic noise on a daily basis  
from an array of seismometers in Parkfield Ca.

Typical Daily Correlation between two of the stations:



Very hard to interpret.

The correlations  
are about 80% converged.

No clear "arrivals."

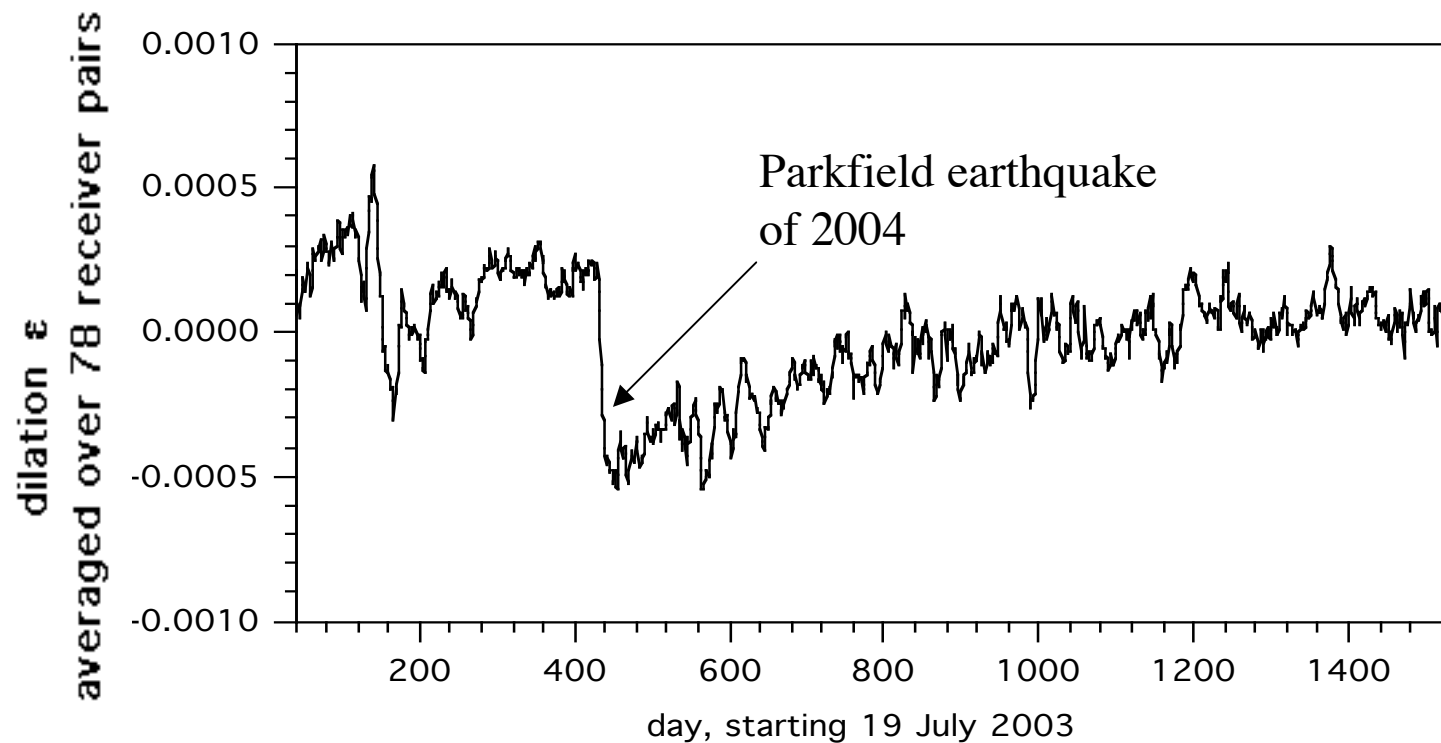
This is G??

They then constructed dilation correlation coefficients

$X(\epsilon)$  (a measure of relative stretch)

( This is a 4th order statistic on the seismic field )

Between the  $C_{ij}$  on different dates and the (year-long) average  $C_{ij}$



Method has been used to predict  
volcano eruptions

In Sum . .

It has been about 12 years now, and the topic is still growing, still hot,  
especially in seismology

Applications in

High resolution seismic velocity maps

Maps of attenuation too?

Monitor changes in a medium

Maps of scattering ?

We still need better understanding  
of the effects of imperfectly diffuse fields.