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**Department of Electrical and Computer Engineering**

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# **Theory and Practice – Making use of the Barkhausen Effect**

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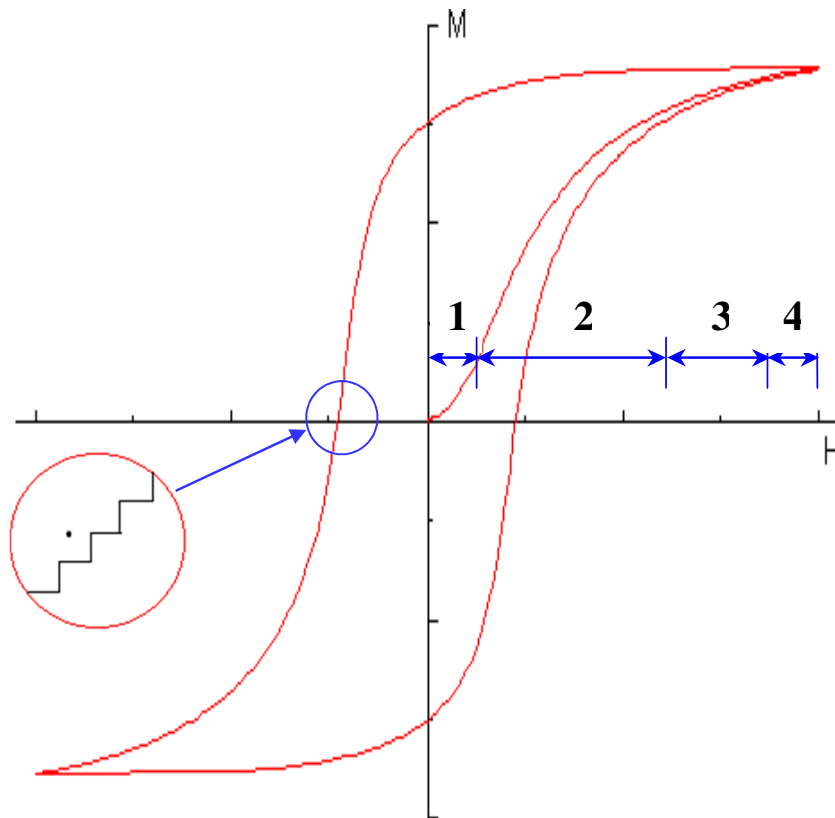
**Workshop on Large Fluctuations and Collective Phenomena  
in Disordered Materials**

**May 18, 2011**

# Summary

- **Stochastic/deterministic Barkhausen model**
  - Equations describing the phenomenon
  - Comparison of theory and experiment
- **Variation of emissions with other factors such as stress**
  - Variation of Barkhausen signal amplitude
  - Theoretical predictions – inverse law
- **Variation of signal with frequency and distance**
  - Possibilities for depth profiling of properties
  - Challenges

## The variation of magnetization with magnetic field looks deceptively simple



- ❑ Yet behind it lies a large number of very complicated interrelated mechanisms which are hard to model
- ❑ A further problem is that magnetic behavior is inhomogeneous and does not “scale” easily as dimensions change
- ❑ The result has been a number of different models which do not fit together very well

## Langevin-Weiss Model

- An array of magnetic moments at temperature  $T$  and in a magnetic field  $H$  will distribute themselves among the available energy states according to the probability distribution  $P$

$$P = P_0 \exp(-\mu_0 m H / k_B T)$$

- Integrating over all possible angles and normalizing gives

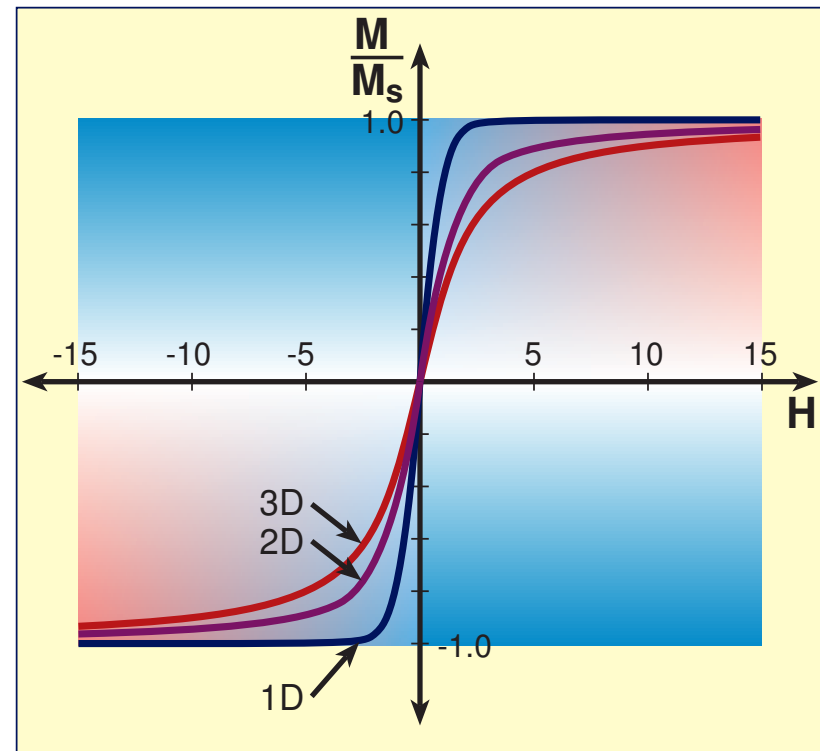
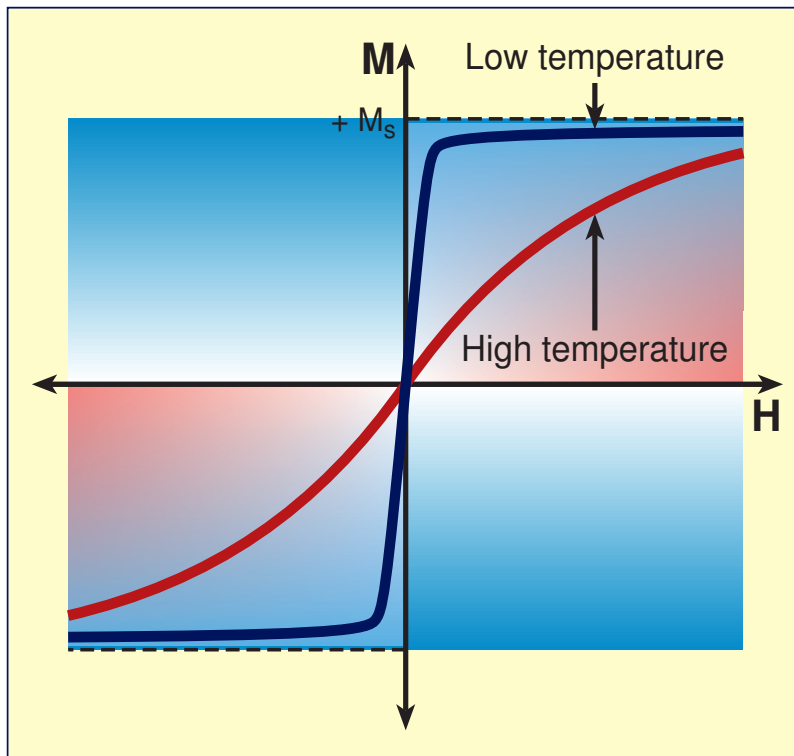
$$M = Nm [\coth(\mu_0 m H / k_B T) - (k_B T / \mu_0 m H)]$$

- Weiss extended this to include internal coupling proportional to the magnetization  $H_e = \alpha M$  with the result

$$M = Nm [\coth(\mu_0 m (H + \alpha M) / k_B T) - (k_B T / \mu_0 m (H + \alpha M))] ]$$

which describes the **anhysteretic** magnetization curve for ferromagnets

## Anhysteretic magnetization curves with different temperatures and different anisotropies



"Superparamagnetic magnetization equation in two dimensions," D.C. Jiles, S.J. Lee, J. Kenkel and K. Metlov.  
*Applied Physics Letters*, 77, 1029, 2000.

## Energy Dissipation and Hysteresis

- Energy is only dissipated by irreversible magnetization processes
- A model with dissipation proportional to change in  $M$  gives good results

$$E_{pin}(M) = \frac{n\langle \mathcal{E}_\pi \rangle}{2m} \int_0^M dM_{irr} = \mu_0 k \int_0^M dM_{irr}$$

- Energy output (change in magnetostatic energy) must equal energy input minus energy dissipated due to any losses such as hysteresis

$$\mu_0 \int M_{irr} dH_e = \mu_0 \int M_{an} dH_e - \mu_0 k \int dM_{irr}$$

## Irreversible Changes in Magnetization

- Differentiating the energy equation lead to a differential for the rate of change of magnetization

$$\frac{dM_{irr}}{dH_e} = \frac{1}{k} (M_{an} - M_{irr})$$

- This can be written in the form of a differential with respect to the applied field H

$$\frac{dM_{irr}}{dH} = \frac{1}{k - \alpha(M_{an} - M_{irr})} (M_{an} - M_{irr})$$



## Isotropic Model of Hysteresis

- The irreversible component of magnetization varies according to the differential equation

$$\frac{dM_{irr}}{dH} = \frac{1}{k - \alpha(M_{an} - M_{irr})} (M_{an} - M_{irr})$$

- The reversible component of magnetization varies as

$$\frac{dM_{rev}}{dH} = c \left( \frac{dM_{an}}{dH} - \frac{dM_{irr}}{dH} \right)$$

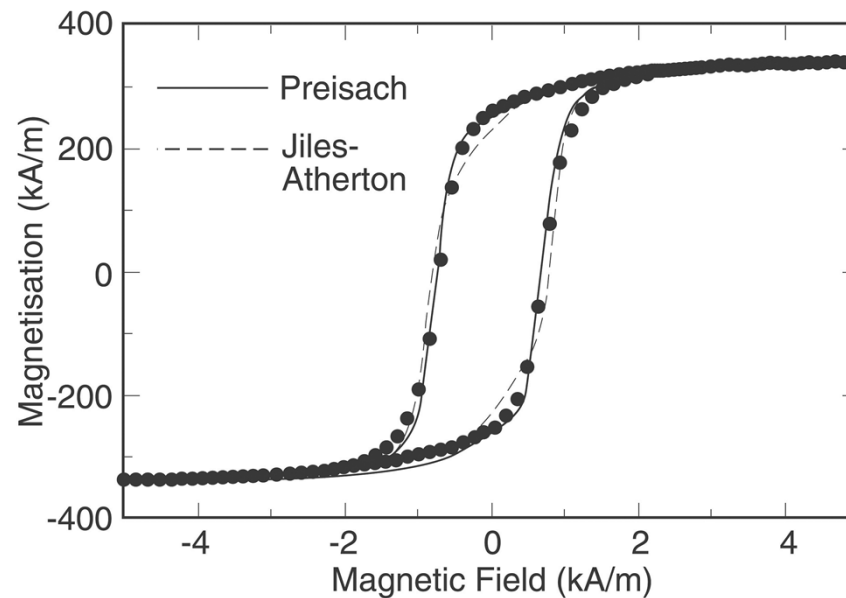
- Therefore hysteresis can be represented in terms of  $M_s$ ,  $a$ ,  $\alpha$ ,  $k$ , and  $c$  using the equation

$$\frac{dM}{dH} = \frac{1}{(1+c)} \frac{(M_{an} - M)}{k - \alpha(M_{an} - M)} + \frac{c}{(1+c)} \frac{dM_{an}}{dH}$$

"Determination of theoretical parameters for modelling bulk magnetic hysteresis properties using the theory of ferromagnetic hysteresis," D.C. Jiles, J.B. Thoeke and M.K. Devine. *IEEE Trans. Mag.*



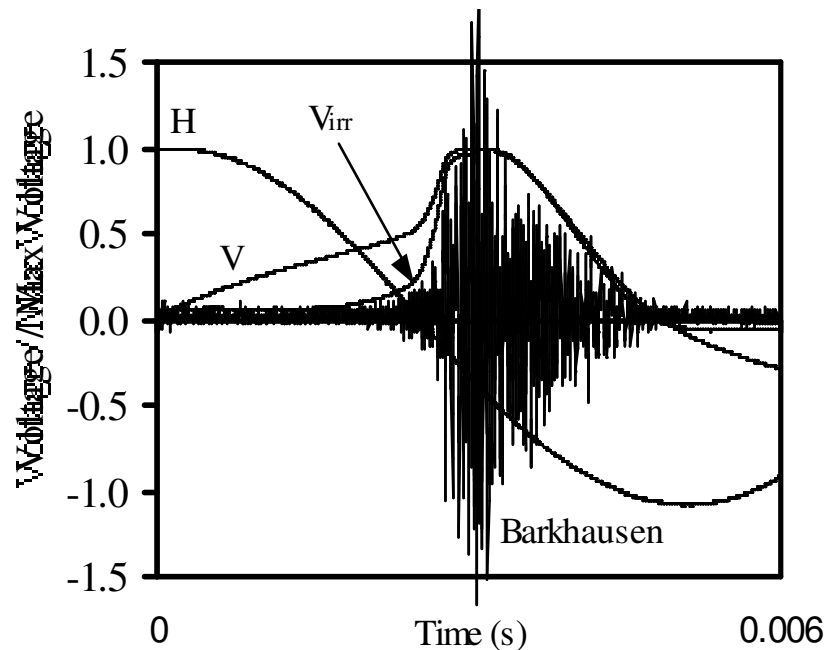
## Comparison of Model and Experimental Measurement



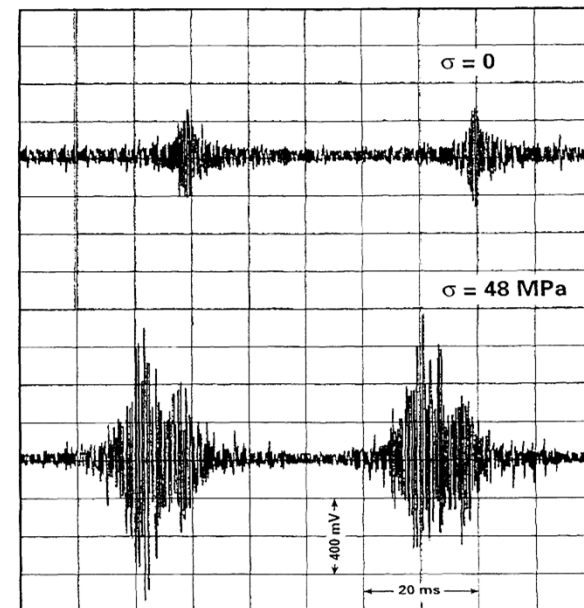
- **Comparison of measured and modelled hysteresis curves of cobalt modified gamma iron oxide material**

After P. Andrei and A. Stancu, "Hysteresis in particulate recording media. Experiment and simulation with Preisach and Jiles-Atherton models," *Journal of Magnetism and Magnetic Materials*, **206**, 160-164, 1999.

## Other magnetic properties - Barkhausen effect



Barkhausen signal shown as voltage against time. The time dependence of magnetic field  $H$ , voltage in flux coil  $V$ , and envelope of irreversible component of voltage in flux coil  $V_{irr}$  are also shown on the same time scale.



Changes in Barkhausen emissions in carbon steel as a result of applied tensile stress.

## Stochastic Process Barkhausen model

Model assumes Barkhausen activity is proportional to the rate of irreversible change in magnetization

$$\frac{dM_{BE}}{dt} \propto \dot{M}_{irr} = \chi'_{irr} \dot{H} \quad M_{BE} = N \langle M_{disc} \rangle$$

$\langle M_{disc} \rangle$ : Average discontinuous change in magnetization due to Barkhausen jump

**Avalanches:** Number of Barkhausen events  $N(t_n)$  in a given time interval is correlated with number of events in the previous time interval  $t_{n-1}$

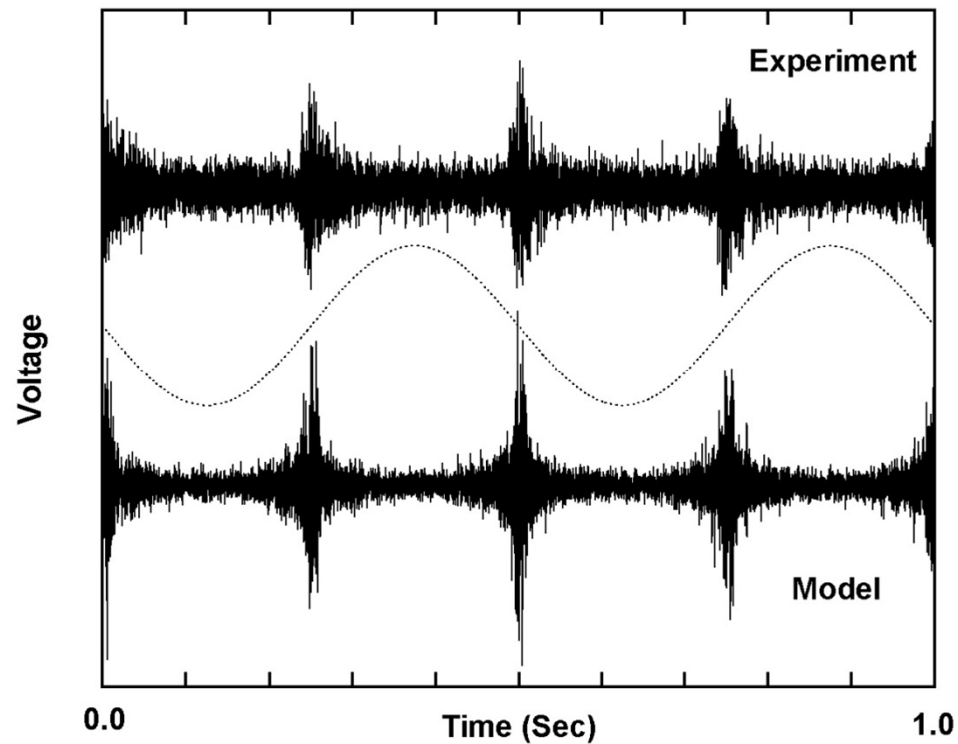
$$N(t_n) = N(t_{n-1}) + \Delta N(t_{n-1})$$

$$\Delta N(t_{n-1}) = \delta_{rand} \sqrt{N(t_{n-1})}$$

$\delta_{rand}$ : Random number  
(-1.47 <  $\delta_{rand}$  < 1.47)

$$\dot{M}_{BE}(t_n) = \langle M_{disc} \rangle \chi'_{irr} \dot{H} \left[ N'(t_{n-1}) + \delta_{rand} \sqrt{N'(t_{n-1})} \right]$$

## Comparison of model with experiment



## Other Factors – Effects of Stress

- Applied stress can be treated in most respects like an effective magnetic field which changes the anisotropy of the material

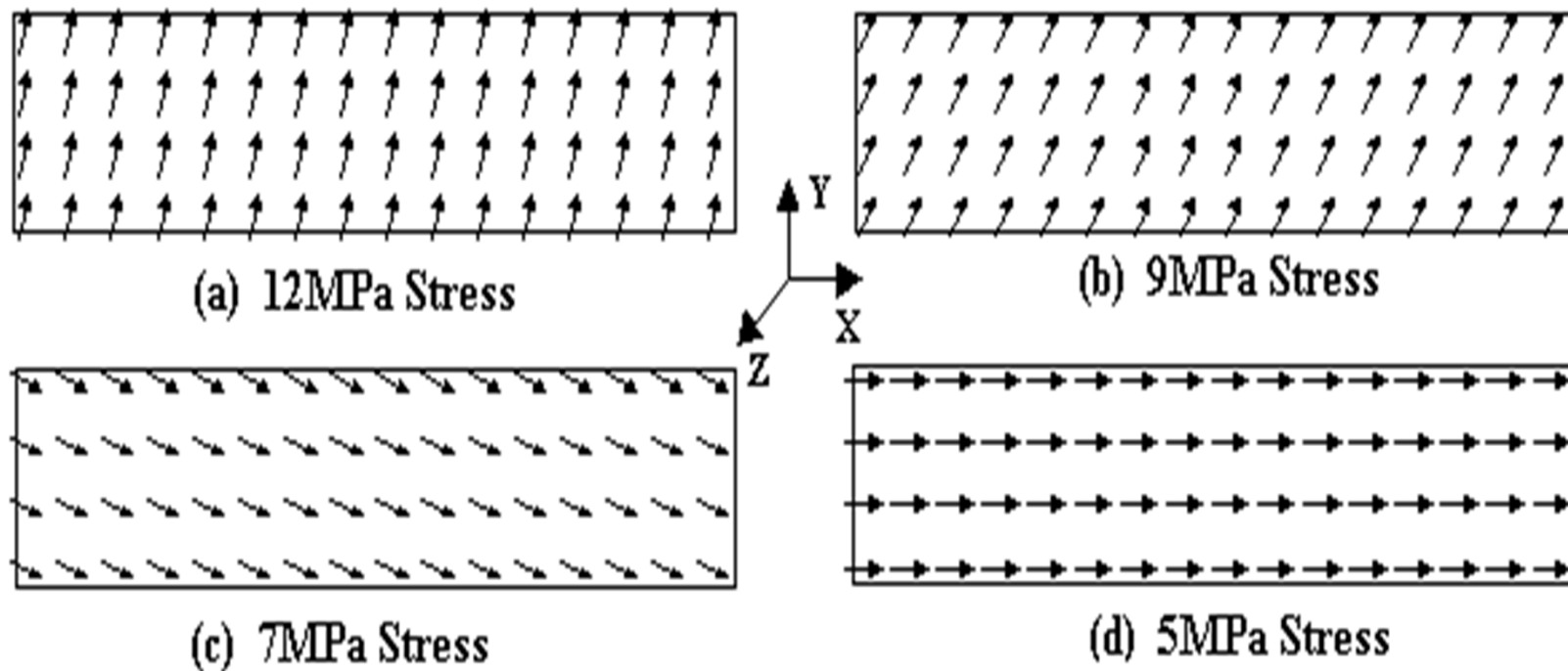
$$H_{\sigma} = \frac{3 \sigma}{2 \mu_0} \left( \frac{\partial \lambda}{\partial M} \right)_T$$

- Later this was extended to cover the case of a uniaxial stress at an arbitrary direction to the applied magnetic field

$$H_{\sigma}(\theta) = \frac{3 \sigma}{2 \mu_0} (\cos^2 \theta - \nu \sin^2 \theta) \left( \frac{\partial \lambda}{\partial M} \right)_T$$

- $\sigma$  is the stress,  $\theta$  is the angle between the stress axis and the direction of  $H_{\sigma}$ , and  $\nu$  is Poisson's ratio.

## Modelled Effects of Stress on Magnetic Moment Orientation



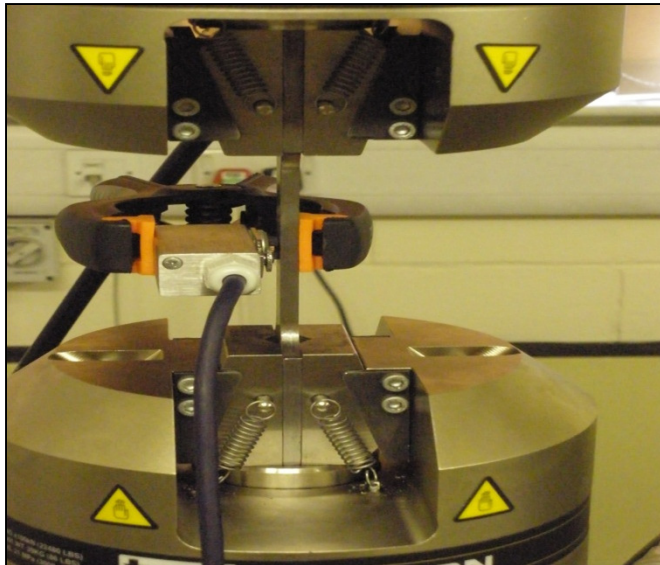
Positive magnetostriction stress applied along the y axis



# Detection of Stress

Specimens calibration tested under load in Instron Tensile Test machine

BN sensor in the middle of gauge length of sample



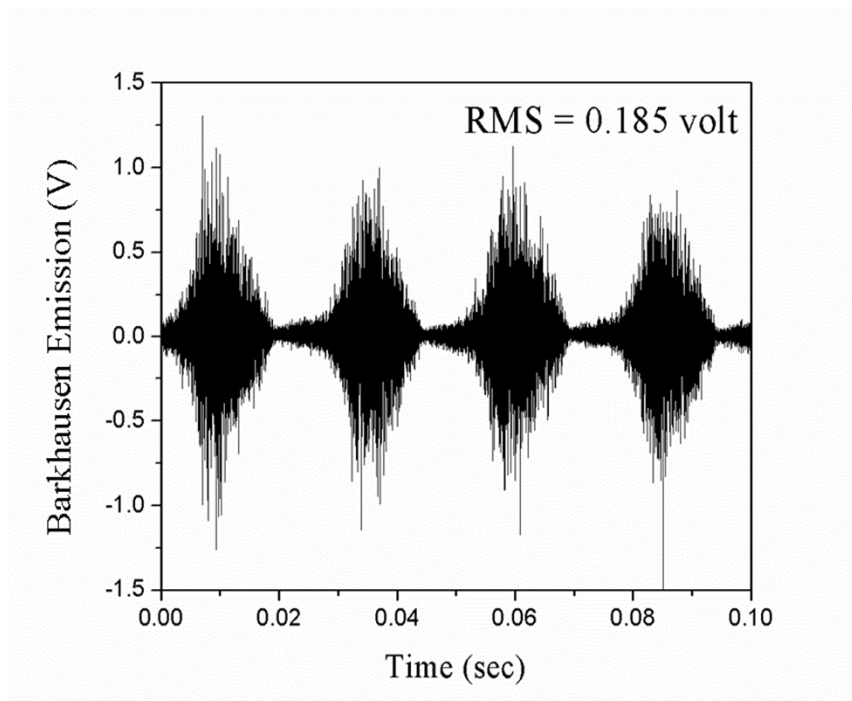
## Typical Magnetising and Analysis Conditions :

**Magnetising frequency :**  $f = 300 \text{ Hz}$   
**Magnetising voltage :**  $V = 12 \text{ volt}$   
**No of bursts :** 20  
**Analysing bandwidth frequency :**  $f_a = 20\text{-}1250 \text{ KHz}$   
(about nearest  $100 \mu\text{m}$ )  
**Smoothing parameter:**  $sp = 100$   
**Sampling frequency:**  $f_s = 2.5 \text{ MHz}$

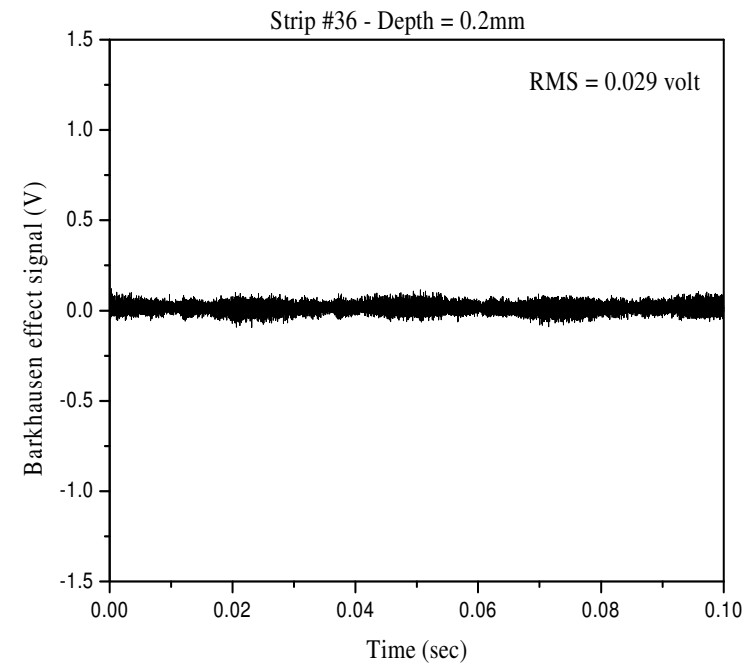


# Measured Barkhausen Results

Core



Surface



## Effects of Stress on Barkhausen Emissions

Applied stress causes a change in anisotropy energy of magnetic moments

$$E_{\sigma} = -\frac{3}{2} \sigma \lambda (\cos^2 \theta - \nu \sin^2 \theta)$$

This can be expressed as an equivalent field

$$H_{\sigma} = -\frac{1}{\mu_0} \frac{\partial E_{\sigma}}{\partial M} = \frac{3}{2} \frac{\sigma}{\mu_0} \left( \frac{\partial \lambda}{\partial M} \right)_{\sigma}$$

The total field is then the sum of magnetic field, exchange field and “stress-equivalent field”  $H_{\sigma}$

$$H_e = H + H_{\sigma} + \alpha M = H + \left( \frac{3\lambda_s \sigma}{\mu_0 M_s^2} \right) M + \alpha M$$

The magnetization is then an anisotropic function of this total field

$$M_{an}(H, \sigma) = M_s \left[ \coth \left( \frac{H + H_{\sigma} + \alpha M}{a} \right) - \frac{a}{H + H_{\sigma} + \alpha M} \right]$$

## Effects of Stress on Barkhausen Emissions

$$\chi'(\sigma) = \frac{dM_{an}}{dH}(H, \sigma) \cong M_s \left[ \frac{1}{3a - (\alpha_\sigma + \alpha)M_s} \right] = M_s \left[ \frac{1}{3a - \left( \frac{3\lambda_s \sigma}{\mu_o M_s^2} + \alpha \right) M_s} \right]$$

This predicts how the rate of change of magnetization with field depends on stress and so can be used to calculate stress. However there is a much easier way...

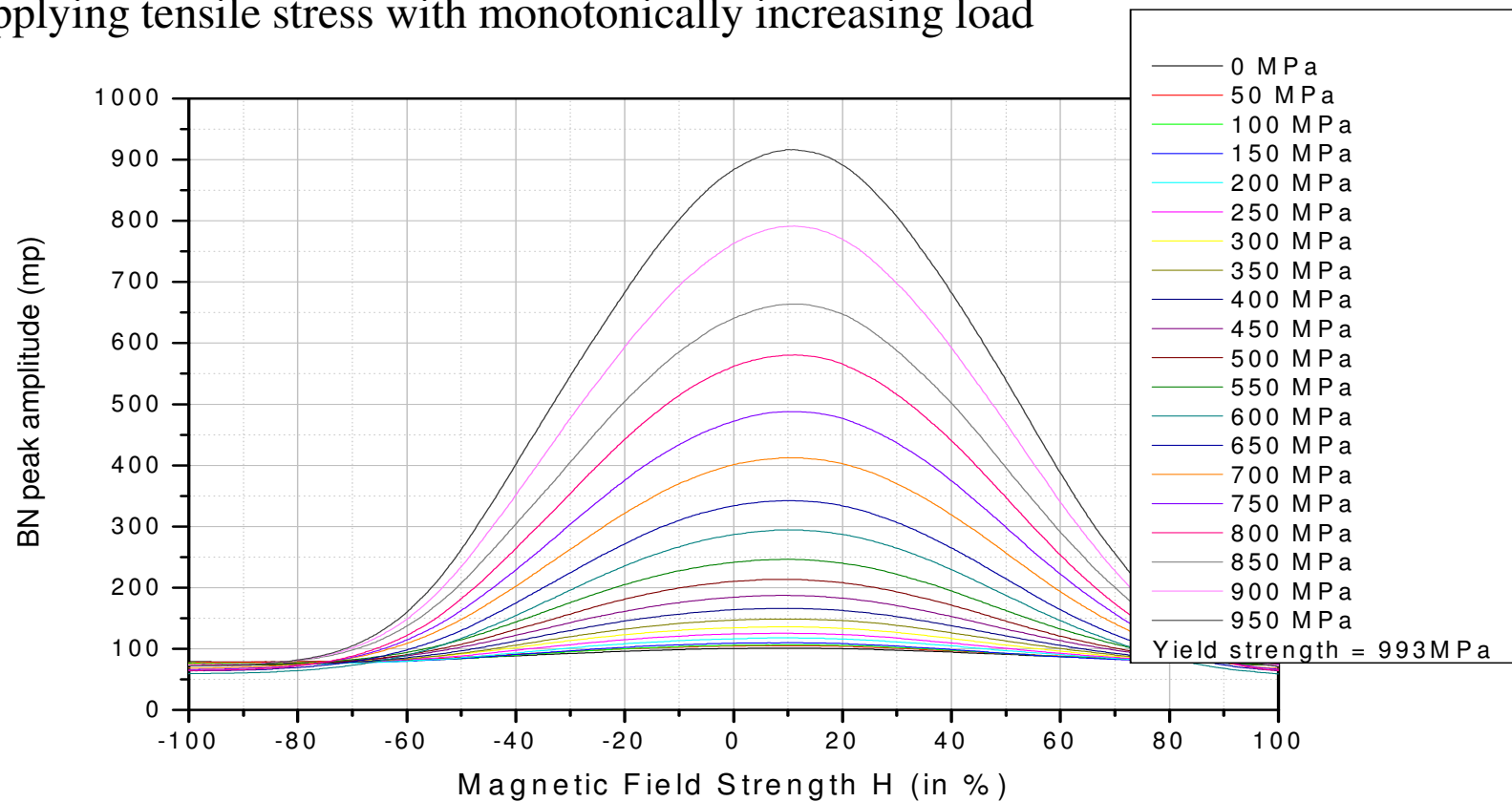
$$\frac{1}{\chi'(\sigma)} = \frac{3a - (\alpha + \alpha_\sigma)M_s}{M_s} = \frac{3a - \left( \alpha + \frac{3\lambda_s \sigma}{\mu_o M_s^2} \right) M_s}{M_s}$$

$$\frac{1}{\chi'(\sigma)} = \frac{1}{\chi'(0)} - \frac{3\lambda_s \sigma}{\mu_o M_s^2}$$

which predicts a straight line graph of  $1/\chi'$  against  $\sigma$

## Detection of Tensile Stress - test results

Applying tensile stress with monotonically increasing load



## Dependence of Barkhausen emissions on stress

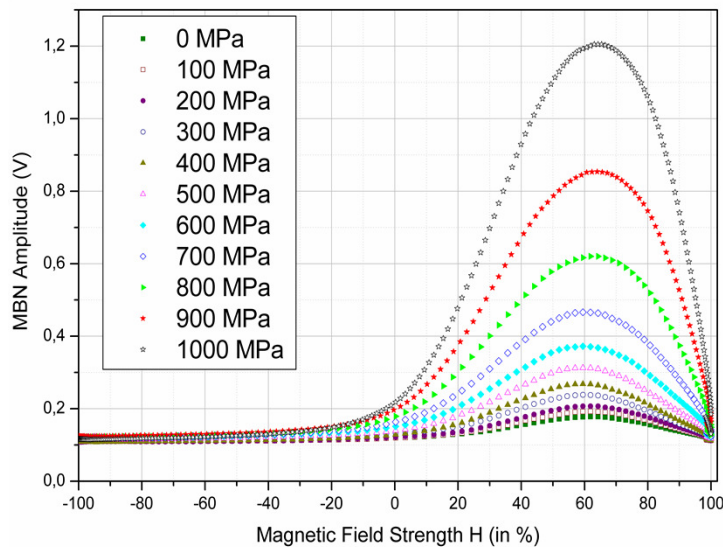


Fig. 1. Envelope curves of the rectified MBN bursts for carburized SAE 9310 specimen for different amplitudes of applied stress using a tensile test machine (residual stress value from XRD measured before stress application was -805MPa).

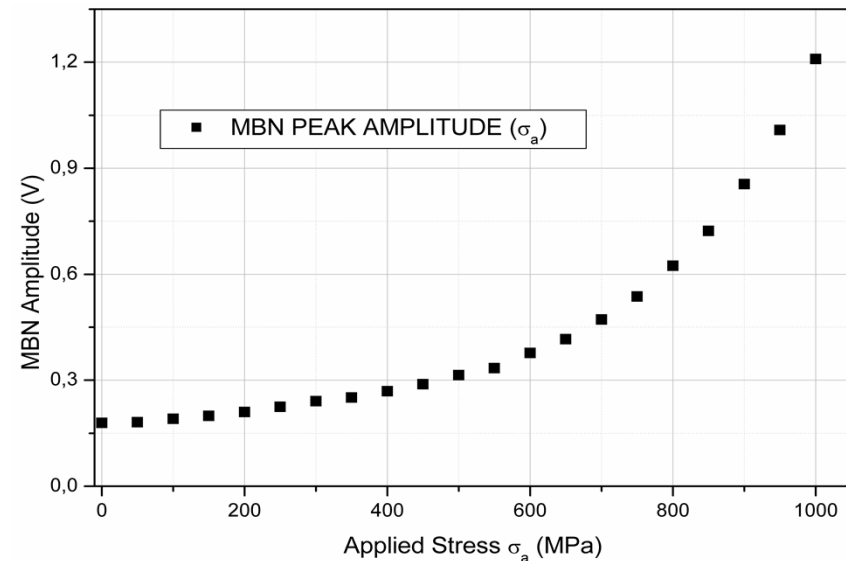


Fig. 2. MBN Peak Amplitude for carburized SAE 9310 specimen as a function of applied stress.

## Effects of Stress on Differential Susceptibility and Barkhausen Emissions

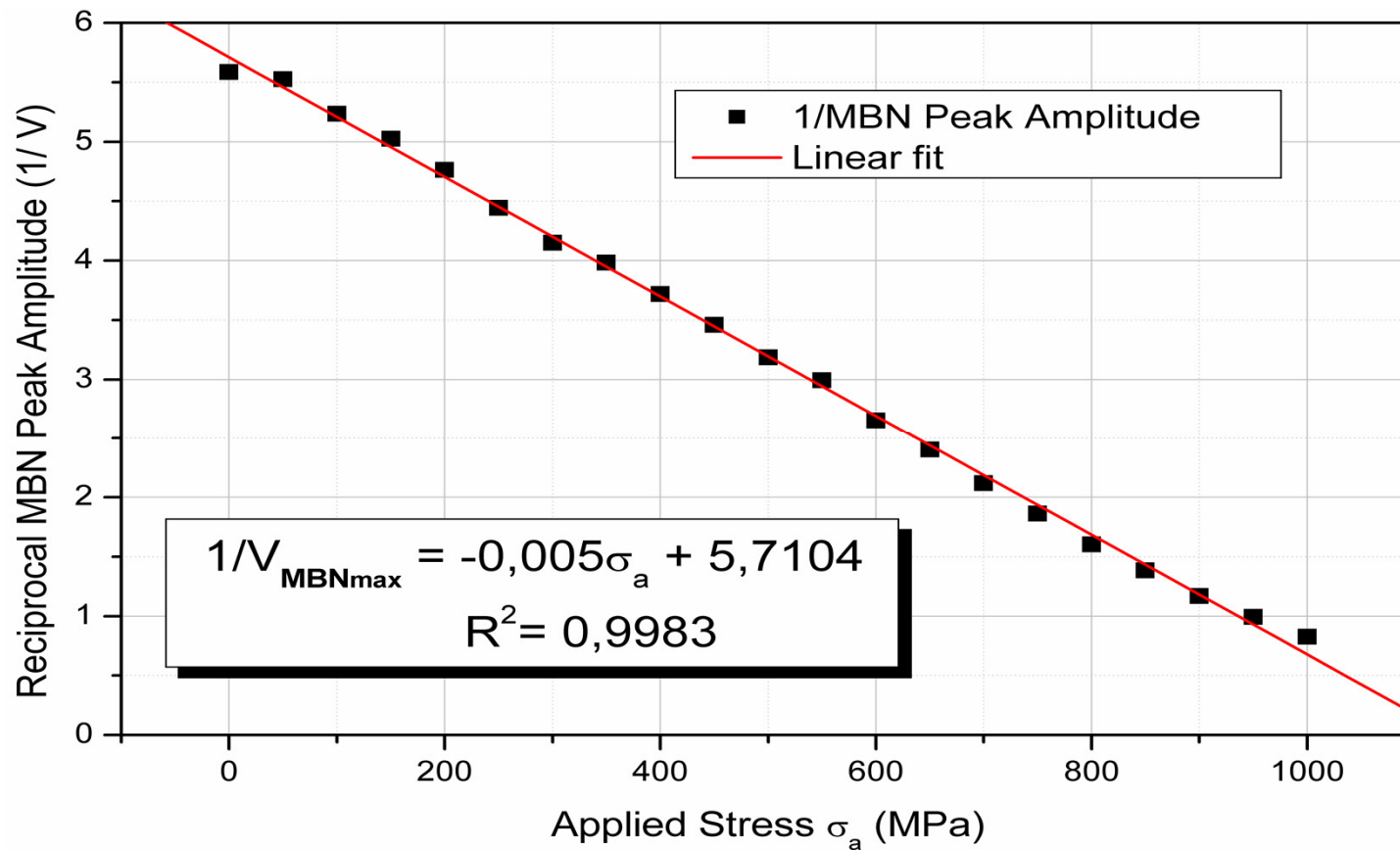
- **Barkhausen voltage  $V_{MBE}$  is known from previous work\* to be proportional to the differential susceptibility**
- **Therefore, a similar linear expression should hold for reciprocal Barkhausen voltage as a function of stress as for the reciprocal differential susceptibility**

$$\frac{1}{V_{MBE}(\sigma)} = \frac{1}{V_{MBE}(0)} - \frac{3b'\sigma}{\mu_o}$$

- **This suggests that the most useful calibration curve for Barkhausen effect as a function of stress is the reciprocal plot**



## Comparison of Model equation with Measurements





## Conclusions

- **Magnetic properties, including permeability and Barkhausen effect depend on other external factors such as stress.**
- **A phenomenological/stochastic model has been developed to describe Barkhausen and these effects**
- **Development of models is essential for understanding and interpretation of measurement results and for predicting changes, such as with stress.**
- **Measurements of these properties can be used for determination of stress and even its variation with depth.**



## Depth Profiling using Barkhausen Effect

There are three key equations

- The change in flux density that propagates to the surface is

$$B_{meas}(0, x_{max}, \omega_1, \omega_2) = \int_{\omega_1}^{\omega_2} \int_0^{x_{max}} B_{origin}(x, \omega) \exp\left(-\frac{x}{\delta}\right) dx d\omega$$

- The voltage measured in an induction coil on the surface is

$$V_{meas}(0, x_{max}, \omega_1, \omega_2) = NA \int_{\omega_1}^{\omega_2} \int_0^{x_{max}} \frac{dB_{origin}(x, \omega, \sigma)}{dH} \frac{dH(x)}{dt} \exp\left(-\frac{x}{\delta}\right) dx d\omega$$

- The component of the measured signal coming from a particular depth is

$$V'(x, \omega_1, \omega_2) = \left( \frac{-4}{x^3 \mu_o \mu_r \sigma} \frac{d}{d\omega} V_{meas}(0, x_{max}, \omega_1, \omega_2) \right)$$

## Comparison with X-ray results

	Ground side surface	Unground initial surface	TOP	BOTTOM
<b>BN PEAK VALUE</b>	939,9	106,0	1416,6	1035,4
<b>BN PEAK POSITION</b>	59,15	48,6	45,4	46,2

